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WT-1401

## OPERATION PLUMBBOB-PROJECT 1.1

BASIC AIRBLAST PHENOMENA (U)
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## FOREWORD

This report presents the final results of one of the 46 projects comprising the military-effect programs of Operation Plumbbob, which included 24 test detonations at the Nevada Test Site in 1957.

For overall Plumbbob military-effects information, the reader is referred to the "Summary Report of the Director, DOD Test Group (Programs 1-9)," ITR-1445, which includes: (1) a description of each detonation, including yield, zero-point location and environment, type of device, ambient atmospheric conditions, etc.; (2) a discussion of project results; (3) a summary of the objectives and results of each project• ${ }^{+}$nr (4) a listing of project reports for the militaryeffect programs.


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Chapter 1
BASIC AIRBLAST PHENOMENA, SHOT PRISCILLA
1.1 INTRODUCTION

Basic measurements of airblast resulting from the detonation of nuclear weapons are necessary to evaluate blast damage to targets of military sign’ ance. As a result of measurements obtained during past tests of nuclear explosions, the presure-time behavior of shock waves producing overpressures in the 5 - to $100-$ psi range is fairly well known. To determine the considerations involved in designing structures to withstand damage at higher pressure limits than were set previously, there has developed a need for investigations of basic phenomena at pressure levels in excess of 100 psi .
1.1.1 Objectives. The primary objective of Project 1.1 was to obtain data on overpressure versus time and dynamic pressure versus time as a function of distance during Shot Priscilla. The pressure range of interest in this shot was from 5 to 1,000 psi for overpressure and from 1 to 650 psi for dynamic pressure.

Additiona. objectives were to: (1) obtain free-field blast measurements at specific locations as required by varıous organizations conducting equipment or structure tests during Shot Priscilla and (2) continue the evaluation of modifications in gage designs, instrument components, and measuremen! techniques.
1.1.2 Background. Basic blast measurements from atomic detonations are needed to properly er luate blast damage to various types of structures of military significance. The blast measurements made during Operation Upshot-Knothole (References 1 through 3) were not sufficient to fully define the parameters most significant in causing damage. In particular, the occurrence of precursor phenomena and distorted waveforms made uncertain the relations betwe en overpressure and dynamic pressure. During Operation Teapot (References 4 and 5) as well as in other operations, considerable information was gathered in the low- and intermediate-overpressure regions for devices of various yields. Within these regions from 5 to 100 psi , overpressure measurements are comprehensive. Although some uncertainty still exists concerning the relation between overpressure and dynamic pressure in the precursor zone, sufficient data was obtained during Operation Teapot (References 4 and 5) from dynamic-pressure measurements to permit predictions on an empirical basis. The test results of Operation Teapot indicate that the factors influencing precursor formation and flow characteristics behind blast waves are surface conditions, height of burst (HOB), and yield of the device.

Recently, emphasis has been placed on the design of structures capable of resisting blast pressures in excess of 100 psi . The requirement for an accurate understanding of the nature of blast waves of these intensities and the desire for accurate knowledge of the pressures acting on the test structures have necessitated measurements at pressures up to $1,000 \mathrm{psi}$. Plumbbob is the first operation where such measurements have been made.

In past tests, asymmetries in blast pressure contours have been noted (Reference 4), that is, a pressure indicated at a given distance along one radius from ground zero may not be found at
is dust-laden, compressibility corrections are also necessary for the air phase portion of the wave to indicate free-stream conditions. Hence, a more realistic comparison with the ideal values would be to use appropriate corrections to the as-read data.

It was realized that agencies taking dynamic pressure measurements had different types of gages and used different notations for the response parameters of the gage as well as different correction techniques for angle of flow and Mach compressibility to yield free-stream dynamic pressure. Correct comparison of this data was difficult.

After the completion of the field test phase following Operation Plumbbob, a conference was called (Reference 9) involving several agencies participating in nuclear field tests to resolve existing problems concerning the measurement and proper interpretation of dynamic pressure, particularly when the blast wave was dust-laden. The agencies represented were Headquarters Armed Forces Special Weapons Project (AFSWP) (now Defense Atomic Support Agency, DASA) Field Command, AFSWP; Naval Ordnance Laboratory (NOL); Stanford Research Institute (SRI); BRL; and Sandia Corporation (SC). Agreement was reached at this meeting to standardize the nomenclature for reporting of dynamic pressure data and to also apply the appropriate correction factors for each type gage used. The nomenclature agreed upon and response of each gage are given in Table 1.3 and Figure 1.7, respectively.

The correction factor for Mach compressibility and corrections for angle of fluw, where applicable, are presented in this report for the BRL q-gages only. Further information about the treatment of data concerning the remaining gages can be obtained from Reference 10.

The as-read value from the BRL q-gage when no dust is present behind the shock wave is given by:
or

$$
\begin{equation*}
q_{c}^{\prime}=\left(P_{t}-P_{S}\right)^{\prime} \text { for } M<1 \tag{1.1a}
\end{equation*}
$$

$$
\begin{equation*}
q_{c}^{\prime}=\left(P_{p}-P_{s}\right)^{\prime} \text { for } M>1 \tag{1.1b}
\end{equation*}
$$

For the clean air case, the free-stream dynamic pressure is defined by:

$$
\begin{equation*}
q=\frac{1}{2} \rho u^{2} \tag{1.2}
\end{equation*}
$$

The relationships for determining the free-stream dynamic pressure behind a clean blast wave from the measurements taken expressed in terms of isentropic flow and normal shock wave equations (Reference 11) are:
For $M<1$

$$
\begin{equation*}
\frac{P_{t}^{\prime}}{\mathbf{P}_{s}^{\prime}}=\left(1+\frac{\gamma-1}{2} \mathbf{M}^{2}\right)^{\gamma / \gamma-1} \tag{1.3}
\end{equation*}
$$

The value of $q$ can be related to the static pressure $P_{S}$ and to the Mach number $M$ in accordance with the following relations:

$$
\begin{equation*}
q=\frac{1}{2} \rho u^{2}=\frac{\frac{1}{2} \rho u^{2}}{C^{2}} \quad \frac{\gamma P_{S}}{\rho}=\frac{\gamma M^{2} P_{S}}{2} \tag{1.4}
\end{equation*}
$$

Where: $\quad C^{2}=\frac{\gamma P_{S}}{\rho}$
To correct the $q_{c}^{\prime}$ value of Equation 1.1a use is made of Equations 1.3 and 1.4, yielding the following:

$$
\begin{equation*}
q=\frac{\left(P_{t}-P_{S}\right)^{\prime}}{\frac{2}{\gamma M^{2}}\left[\left(1+\frac{\gamma-1}{2} M^{2}\right)^{\gamma / \gamma-1}-1\right]} \tag{1.5a}
\end{equation*}
$$

or:


Equations 1.5 a and 1.5 b are known as the pitot tube equations. Use of the latter two equations becomes superfluous, since it would suffice to use Equation 1.3 to obtain the Mach number and by Equation 1.4, using $\gamma=1.4$, the $q$ or free-stream dynamic pressure is derived directly. Another advantag 2 for using Equation 1.4 is that, for supersonic flow, it is necessary to correct the total head pito! pressure $P_{p}$ to the free-stream total pressure $P_{t}$ prior to use of Equations 1.5 a or 1.5 b .

Hence: For $M>1$

$$
\begin{equation*}
\frac{\mathbf{P}_{\mathrm{t}}}{\mathrm{P}_{\mathrm{p}}},=\left[\left(\frac{2 \gamma}{\gamma+1} \quad \mathbf{M}^{2}-\frac{\gamma-1}{\gamma+1}\right)\right]^{\frac{1}{\gamma-1}}\left[\frac{1+\frac{\gamma-1}{2} \mathrm{M}^{2}}{\frac{\gamma+1}{2} \mathbf{M}^{2}}\right]^{\gamma /(\gamma-1)} \tag{1.6}
\end{equation*}
$$

anld:

$$
\begin{equation*}
\frac{\mathbf{P}_{\mathbf{p}}^{\prime}}{\mathbf{P}_{s^{\prime}}^{\prime}}=\left[\frac{\left(\frac{\gamma+1}{2} \mathbf{M}^{2}\right)^{\gamma}}{\left(\frac{2 \gamma}{\gamma+1} \mathbf{M}^{2}-\frac{\gamma-1}{\gamma+1}\right)}\right]^{1 /(\gamma-1)} \tag{1.7}
\end{equation*}
$$

Obviously, Equation 1.6 is not necessary when use is nisde of Equation 1.4.
In Summary, for $M<1$, Equation 1.3 along with Equation 1.4 yields the free-stream dynamic pressure, and Equation 1.7 combined with Equation 1.4 yields the free-stream dynamic pressure for $M>1$.

The above considerations are for the clean blast wave. When the blast wave is dust-laden, the free-stream dynamic pressure is defined by the combination of the airflow and momentum flux of dust, i.e.,

$$
\begin{equation*}
q_{c}^{* \prime}=q+\phi_{d} \tag{1.9}
\end{equation*}
$$

The response of the BRL q-gage is given by:

$$
\begin{equation*}
\mathrm{q}_{\mathrm{c}}^{* \prime}=\left(\mathrm{p}_{\mathrm{t}}^{* \prime}-\mathrm{p}_{\mathrm{s}}\right) \text { for } \mathrm{M}<1 \tag{1.10a}
\end{equation*}
$$

or:

$$
\begin{equation*}
q_{c}^{*^{\prime}}=\left(P_{p}^{* \prime}-P_{S}\right) \text { for } M>1 \tag{1.10b}
\end{equation*}
$$

This may also be expressed by the following equation:

$$
\begin{equation*}
q_{c}^{* \prime}=\left(q_{c}+n \phi_{d}\right)^{\prime} \tag{1.11}
\end{equation*}
$$

Preliminary investigations of the dust registry coefficient $n$ for the BRL q-gage shows the value to be small (Reference 12) lying between 0.12 and 0.21 .

The problem of obtaining free-stream dynamic pressure in the presence of dust, therefore, is to measure separately the contributions to the dynamic pressure resulting from air alone and that resulting from dust. To realize this condition, at least two different gages are required, with known registry coefficients for both. With these two values known, the momentum flux could be determined and then from Equation 1.11, the air portion $q_{c}$ would be known. Expressing the above mathematically would yield the following:



The difficuitise encountered and experience gained during the past operations led to many modifications of the gages for use during Operation Dlumbbob. Figures 1.9 through 1.11 show the $P_{t}$-gage, $q$-gage, and the new prototype $q$-gage.

Pressure versus Time Gages. Some minor modifications were made to the $\mathbf{P}_{t^{-}}$ gages. A new thermal link was designed in which the metal creep in the thermally sensitive solder is minimized.

The photo-initiation circuit was changed to keep the initiation circuit separate from any timing circuit. The initiation circuit was repackaged with circuit components potted. All changes were made to make the $P_{t}$-gage an even more reliable and accurate instrument while maintaining its portability and ease of installation and recovery.

In addition, the gages at the stations close to ground zero were modified to be initiated by a timing signal. The hard-wire initiation was used on one of the two gages on the first five stations on the blast line. The hard-wire-initiated gages are listed in Table 1.6 with a letter "B" after the station number. The signal for the gages was produced by the closure of a standard Edgerton, Germeshausen and Grier (EG\&G) timing relay at -5 seconds and was used to close the electrically latching relay. The operation of the -5 -second relay was duplicated by a second relay at -1 second.

The thermal initiator used during the $P_{t}$-gage during Operation Plumbbob operated similarly to the initiator used during Operation Redwing. However, to prevent premature initiation caused by creep in the solder material, the overlapped portion of the thermally sensitive link was modified so as to require actual separation of the two portions before initiation of the gage could take place.

The photo-initiation system, which was actuated by the sharp transient character of incident light from the detonation, employed a cadmium-sulfide cell, a transistorized amplifier, and a sensitive relay. In field use, a density-of-3 neutral filter was placed directly over the photocell to reduce the quiescent current and to eliminate pre-initiation caused by random pulses of light.

A schematic diagram of the complete photo-initiation circuit of the pressure-time gage is shown in Figure 1.12. To prepare the gage for recording, the activating microswitch SWI, which completes the ground return circuit and which is normally open, was closed. The amplifier circuit is an ordinary grounded emitter utilizing a CK722 PNP transistor. Prior to initiation, the transistor is biased nearly to cutoff by resistor $\mathbf{R}_{\mathrm{i}}$, which is connected to the positive terminal of the 45 -volt battery. At zero time, light impinging upon the photocell reduces the resistance by a factor oí approximately $10^{6}$. As a result, a large negative pulse appears at Point A. This negative pulse, coupled through Capacitor C1 to the base of the transistor, causes the transistor to conduct heavily. This current surge closes the sensitive relay ( $\mathrm{RL}-1$ ), which is in series with the transistor collector and the negative side of the battery.

A pair of normally open contacts on RL-1 were used to latch the relay in the closed position by placing the battery voltage directly across the relay coil. The increased current, 5.6 ma , through the relay insured continuous contact closure, despite blast and ground shocks. The other set of relay contacts was used to apply voltage from the 8 -volt mercury-cell source to the drive motor. This started the recording cycle, and the turntable continued to rotate a predetermined number of revolutions until the star cam opened the normally closed microswitch (SWI), which disconnected the motor and the transistor from their respective voltage supplies.

For protection against dirt and moisture, a portion of the amplifier circuit was potted in an epoxy resin.

The basic component common to the self-recording gages is a pressure-sensing capsule. Fundamentally, the new capsules used during Operation Plumbbob were similar to those used previously. Capsules with ranges from 0 to 5 psi up to 0 to 1,000 psi were used in the selfrecording gar,es. The capsules with ranges from 0 to 800 psi and 0 to $1,000 \mathrm{psi}$ were developed specifically ior use during Operation Plumbbob.

The cha acteristics of the individual pressure capsule were known from the calibration curves supplied by the manufacturer. These curves showed the deflection of the stylus as a function of
the pressure. Laboratory calibration of a group of capsules selected at random showed the manufacturer's calibration curves to be sufficiently reliable to be used without further check.

The heart of the timebase in the pressure-time and dynamic pressure gages is the $A . W$. Haydon, Series 5600, chronometrically governed, dc motor (Reference 14).

A study was undertaken in BRL to determine the reliability of the motor under field conditions. Tests were performed to obtain turntable behavior, as a function of time, from the instant power was applied until it was removed and to study the effect of shock acceleration as great as 80 g .

The first tests determined the consistency of the length of time required for the motor to reach constant velocity. Oscillograph traces were made showing motor current, turntable rotation, and reference time for free motors and for motors loaded in the manner in which they were to be used. This data was then plotted as turntable rotation versus time (Figure 1.13 and 1.14). For the $3-\mathrm{rpm}$ motors, the data shows that they attained full velocity in approximately 100 msec , exceeded their rated velocity, and then settled down to constant speed. Overshoot and oscillation were decreased by loading, but in no case lasted longer than 400 msec . The difference between the plot of total angular displacement versus time and a linear extrapolation to zero rotation of constant velocity never exceeded 20 msec after plus 50 msec from zero time. The time difference between the point of intersection of the extrapolated line with the $x$-axis and true zero as indicated on the plotted curve was 65 msec , with a standard deviation about this value of 11 msec . As a result of the test of $10-\mathrm{rpm}$ motors, the data shows that the motors attained constant velocity in approximately 400 msec without exceeding the rated velocity or oscillating about this velocity prior to reaching constant speed. For these motors the time difference between the point of intersection of an extrapolated line of constant rotation with the $x$-axis and true zero, as shown on the plotted curve, was about 150 msec .

The second group of tests was made to determine the effect of shock on turntable velocity. Again, oscillograph traces were made of turntable rotation, motor current, and reference time while the gage was subjected to shock in a Barry 150 VD medium-impact-shock machine. The motors were tested under impacts as high as 80 g with a duration of 12 msec . Shork-acceleration tests were run in preference to vibration-acceleration tests, because the fornter more closely duplicated field conditions. The gages withstood impact shock of at least $\overline{\mathrm{a}} 0 \mathrm{~g}$ without appreciably affecting record accuracy. Deviations from the linear operation line did not exceed 5 msec under a shock of 50 g . At higher accelerations, the velocity oscillated after the instant of shock in some cases, and occasionally, the glass recording disk broke. Acceleration shock at 50 g or less, however, did not materially affect motor and turntable velocity.

Dynamic Pressure versus Time Gages. The dynamic-pressure gage was a modified model of the type used during Operations Redwing, Teapot, and Castle. The gages used during Redwing were recovered, cleaned and rebuilt. The gage nose section, which contains the capsules and recording elements, was filed and polished to remove most deep scratches and dents suffered during the tests. The main body of the gage was cleaned by sandblasting and refinished with baked enamel.

The major modification to the gage is the installation of the new-type pressure capsules. The turntable assembly was changed to bring about more nearly constant speed of rotation. This involved mounting the drive spindle on roller bearings. The spindle had previously been a running fit in the turntable housing. Because of the tendency of the spindle to bind, turntable rotation had been uneven. To reduce acceleration effects, the mass of the record disk retainer was reduced, and the diameter of that portion in contact with the glass disk was increased to furnish more support for the disk. A new method of mounting the relays, batteries, and other components was designed to prevent intermittent operation due to acceleration, to provide easier means of checking the gage circuits, and to provide for better protection against the elements.

Gage initiation was accomplished through the use of a photoelectric system backed up by a thermel initiator.

The ${ }_{1}$ hotoelectric system for the $q$-gage was different from that system used on the $P_{i}$-gage. It contait ed a 918 phototube, a sensitive plate relay, an adjustable potentiometer, and a mechanical latch.ng relay. The sensitive relay closes when the light becomes sufficiently intease to

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allow $400 \mu \mathrm{a}$ of current to pass. This closure completed the solenoid circuit of the latching relay, which closes the drive motor circuit. This circuit functioned properly on 23 of the 24 dynamic pressure gages used during Shot Priscilla.

The standard BRL q-gage was modified so as to make it possible to mount this gage on a standard contractor-installed mount. The base flange of the BRL self-recording q-gage is identical with the base flange on the electronic dynamic-pressure gages. The BRL q-gage and a standard contractor-installed mount are shown in Figure 1.15. Another modification was the relocation and redesign of the phototube holder and thermal link assembly. This modification in design completely sealed the electronic-initiation system from moisture.

Design of a New Gage for Dynamic Pressure versus Time. In designing a new dynamic-pressure gage, it was considered that the following features were desirable: (1) ability to be changed from a two-element $P_{s}$ and $P_{t}$ to a direct-reading dynamic-pressure gage, (2) simpler field servicing and installation, and (3) use of standardized pressure-timegage components whe rever possible ( $31 / 2$-inch $P_{t}$-gage recording disks, 3 - and 10 -rpm motors, and $P_{t}$-gage initiation system).

Five gages of an experimental model incorporating these features were built and used during Operation Plumbbob.

In design and construction, the new gage represents a radical departure from the BRL $q$ gage used during previous operations. The old gage consisted of a nose in the form of a pitot tube attached to and collinear with the cylindrical body that housed the battery pack and initiation circuits and that was used to support the tube on the field mount. The overall dimensions of the gage, including the pitot tube nose and body, were approximately 36 inches in length by $3 \frac{3}{4}$ inches in diameter. All of the recording mechanism was contained within the pitot tube section. During tests, it was necessary to orient the axis of the gage body along a line toward the blast. With this orientation, the $3 \frac{3}{4}$-inch gage-body diameter restricted the fresdom of placement of gage components and the size of the recording disk.

The new gage has a cylindrical body oriented at right angles to the nos $\equiv$ of the gage. This body houses the pressure-sensing capsules, drive motor, recording sjsiem, and standard $P_{t^{-}}$ gage, and it accommodates a $3 \frac{1}{2}$-inch $P_{t}$-gage disk. Capsules and motors are interchangeable with $P_{t}$-gage components. Turntable bearings are the same as $P_{t}$ turntable bearings. The photoelectric initiation system uses the same components with different mountings. The thermal initiator is the new, smaller, standardized plunger and link support.

The pitot tube assembly protrudes from the side of the cylindrical body and has a removable nose section. This permits conversion of the gage from a two-pressure, manual-subtractiontype gage to a direct-reading differential-pressure gage by the use of interchangeable gaskets at the junction of the nose and nose cone. In addition, the separable nose permits the use of extension noses when it is desirable to place the pressure inlets farther from the main body to avoid blast-wave reflections that would affect the accuracy of the pressure records obtained.

The main body and nose-cone assembly will be described together, since they are closely related. The main body is machined from a solid billet of duralumin. The outside is finish machined and the inside roughed out. The cavity for the insertion of the nose cone is precision bored. The nose cone is completely machined, with the cylindrical portion finished to 0.005 inch larger than the cavity in the body. To assemble the two pieces, the nose cone is soaked in dry ice and the main body heated to $200^{\circ} \mathrm{F}$ in an oven. This temperature difference of approximately $300^{\circ} F$ is sufficient to allow the nose cone to be easily inserted into the main body. The resulting shrink fit forms a permanent fit between the two pieces, and welding or pin fastening is unnecessary. The finish machining of the interior of the main body and nose cone is completed with the two pieces assembled.

To keep the mass of moving parts to a minimum, the rotating parts, turntable, and recordblank keepers are machined from uural. The turntable bearings are $30-\mathrm{mm}$ angular-contact bearings preloaded to minimize acceleration sensitivity. The motor and motor batteries are located in the lower ead of the main body. The top portion contains the photoelectric initiation system with its battery supply and relays, the thermal initiator and switch, a master cutoff switch, and a timing switch. The timing switch is screw actuated and can be preset to cut off
the operation of the gage after completion of 2 revolutions or as many as 12. All covers and sections of the main body are sealed with 0 -rings to maintain a constant air pressure inside the gage.

In the field, the gage is secured in position by means of a dural flange bolted to the base of the gage. The flange is 6 inches in diameter and $3 / 4$ inch thick. Eight $3 / 8$-inch bolts on a $5 \frac{1 / 4}{}$ inch bolt circle are used to fasten the gages to the field mount.

All laboratory tests of the new q-gage were yerformed in the BRL 24 -inch shock tube. The axis of the pitot tube lay along the axis of the shock tube, so that it was completely submerged in the airflow following the shock front. The gage was tested under a variety of conditions and pressures. The pressures varied between 5 and 30 psi side-on, and tests were conducted with both clean and dust-laden air. The gage was positioned with angles of yaw from $0^{\circ}$ to $40^{\circ}$. Several tests were made with the extension pitot tube. A series of tests at various pressures were fired with the gage ported for direct dynamic pressure. Preliminary investigations of the shock-tube data showed good agreement in all cases between the dynamic pressures indicated and those expected. However, because of the questionable performance of this gage under field conditions, postshot investigations and anaiysis of the shock tube data did not seem to be warranted.
1.2.3 Installation of Gage Stations. Surveys for the blast lines were made by the civilian contractor. At stations where the peak overpressure expected was in excess of 35 psi , gage mounts for the $P_{t}$-gages were contractor installed. At the first six stations, the $P_{t}$-gage mounts were of the type designed to accommodate two of these gages. Mounts for the selfrecording q-gage were contractor installed where the dynamic pressure was expected to ive above 200 psi. These mounts were the AFSWP standard-design 3-foot a-gage tower for use in the high-pressure zones (Figure 1.16). At all other gage stations, the mounts were installed by project personnel. For the q-gage, this installation consisted of inserting the gage mount, pouring Cal-Seal (a quick-setting cement) around the mount, and leveling the mount so that the gage would be parallel to the ground and directed toward ground zero. A 3-foot gage mount with the new q-gage installed is shown in Figure 1.17. Ten-foot and 3-foot tower mounts for the $q$-gage were prefabricated. In all cases, the gage or gages required at a particular station were placed in their mounts and secured by project personnel.
1.2.4 Data Reduction and Presentation. Self-recording gage records scribed on a rotating disk present the record in polar coordinate form, and it is necessary to convert to rectangular coordinates. A Gaertner toolmaker's microscope with a rotating table was used for the conversion. The microscope was modified by the addition of digital read-out heads.

The information from the read-out heads on the microscope is converted to digital form by Telecordex equipment and then punched on IBM cards. These cards, representing readings taken at short intervals throughout the span of the record, together with cards representing calibration steps are used as input data for the EDVAC high-speed digital computer for linearizing both time and pressure values. The program coded for the EDVAC uses a straight-line equation, and the pressure values are calculated from a straight-line interpolation between the various calibration steps. The output data from the EDVAC is then punched on IBM cards, and these cards are fed to an Electronics Associate vari-plotter. The plotter has a $30-$ by 30 -inch plotting table and can plot 6.6 points per inch. The plots of the pressure-time histories can be varied in size. After the points of the pressure-time plots are connected by straight lines, the records are reduced in size by photographic means for inclusion in the final report.

Photographs of the plots of the pressure-time history data are shown at the end of this chapter.
1.2.5 Dynamic Pressure Data Reduction. Similar procedures were used as described in Section 1.2 .4 for reading records from the BRL q-gage disks. Several steps were taken to reduce the data from the records for purposes of calculating the free-stream dynamic pressures.


It should be brought out at the outset that it was hardly possible to separate the dust flow from the airflow for Shot Priscilla and also that appropriate corrections could be applied to the pressure measurements' for yaw and pitch angles of flow. Lack of yaw and pitch gages as well as the unknown dust registry coefficient for the BRL gage precluded any such considerations.

There is reason to believe from the preliminary test for determining the dust registry coefficient (Reference 12) that the influence of dust on the pressure measurements is small. Furthermore, except for very close distances, as the distance from ground zero is increased, the air-to-dust mixture decreases, minimizing the influence of dust on the measurements. Hence, some scatter in the data can be expected.

From past tests, it was determined that the pitch and yaw records in the presence of a precursor indicated that flow became parallel to the gage axis within the first 30 to $\mathbf{5 0} \mathbf{~ m s e c}$. Within this period of time for most of the q-gage records within the precursor zone, the flow is not at its maximum value. It is believed that results derived were not seriously hampered by not using the correction factors for pitch and yaw angles of flow.

The initial step for calculating the free-stream dynamic pressure was to compute the Mach number from the ratio of total to static pressure using either Equacions 1.3 or 1.7. Wherever possible the surface overpressure measurements were substituted for the q-gage side-on pressure measurements. In case this could not be done, then a correction for zero angle of flow was applied in accordance with that given in Reference 13 following the application of the initial step. The latter case required that an iteration process be used for computing a new Mach number. The new Mach number, along with static pressure, was substituted in Equation 1.4 yielding the calculated free-stream dynamic pressure. The data was processed as a function of time, and the BRL high-speed computer was used for the computation.

### 1.3 RESULTS

The performance of the gages utilized for the blast line instrumentation of Project 1.1 and free-field blast instrumentation for other projects was generally good. The failures of some of the self-recording pressure-time gages were due to premature initiation, failure of initiation, and high-acceleration effects (particularly on the q-gages). In spite of the premature initiation and failure of initiation, peak values of pressures from the gages were obtained. It can be assumed that, in the tabulation of results where neither the arrival time nor the positive duration data is given, only peak values were recorded on the gages. The q-gages at distances of 1,650 feet or less from ground zero were affected by ground accelerations, resulting in excessive hash and high-frequency oscillations on the records.

Ninety-seven self-recording gages were used to obtain pressure-time versus distance data for various flow phenomena from a nonclassical blast wave. Table 1.6 lists the gages used for each project and the type of record obtained.

The curves presenting the basic blast parameters versus distance included herein have been fitted by eye.
1.3.1 Blast Line Data, Project 1.1. The plots of the measured maximum values of overpressures as a function of distance for the main blast line and the free-field measurements taken for the various projects are shown in Figure 1.18. The tabulated results of the overpressure measurements are given in Table 1.7, along with the arrival time and positive duration values. When the time of arrival and positive duration are not given in the table, then the pressure indicated is peak only. The pressure-time histories are shown in Figures 1.19 through 1.27, which include all the free-field overpressure measurements taken.

The depression of the overpressure distance curve and waveforms of the pressure-time curves indicates the formation of a precursor as was expected in Shot Priscilla. The free-field measurements for the various projects were in the Frenchman Flat area west of ground zero. Generally, the small variation of the maximum pressure values between these scattered measurements and those of the main blast line indicates that the shock wave was symmetrical. This

is further borne out by the comparison of waveforms at the same distance from ground zero. A striking feature of the similarity between pressure-time records at an equivalent ground range is shown in Figures 1.22 and 1.23. Four $P_{t}$ gages were located at a ground range of 1,040 feet, each gage separated along an arc by about 70 feet, the two end gages separated by about 240 feet. The differences in the waveforms at this ground range are in the high-frequency oscillation present on the record.

Plots of the as-measured difference between the total head pressure and the static pressure de shown in Figure 1.28. The calculated or currected dynamic pressures are shown in Figfre $1 . \AA^{\wedge}$. Both figures show data obtained along the main blast line and for various projects. The tabulations of the maximum pressures for total, static, and dynamic (pressure difference and calculated) are given in Table 1.8 for the main blast line. The remaining tabulations of the same data for various projects are given in Tables 1.9 through 1.16. In the tables, only those values are reported for which pressure-time histories were obtained. The pressuretime plots are shown in Figures 1.30 through 1.50. In these figures the computed Mach number as a function of time is also shown. The st tic pressures reported in the tables and shown in the pressure-time figures are in most cases the ground baffle values unless otherwise noted.

In Figures 1.28 and 1.29 the subscript 10 adjacent to the symbols designating pressure values indicates measurements taken at a 10 -foot elevation. All others are at the 3 -foot elevation. A wide scatter of data occurs in the lower portion of the curves of Figures 1.28 and 1.29. The gage ranges for the total head and static pressure are similar, and because of the small pressure difference measured, a large variation can be expected because of reading errors as well as some oscillation on gage records.

An attempt was made to smooth through the high-frequency oscillations of the corrected dynamic pressure-time curves. The smoothing of the pressure-time curves is indicated by the dashed lines.

Although a record was inscribed on the $q$-gage disk at 850 feet, the wave was so distorted from acceleration effects that any value reported would be questionable. The new q-gage results do not compare favorably with those from the old gage. The maximum values are lower and duration times are shorter. The latter however, could be attributed to the fact that the motor did not achieve maximum speed.

The arrival time measurements from the main blast line are not representative, since in some cases the gages initiated by signal wire did not respond in the manner expected and therefore are misleading. However, by taking all the data obtained from various projects and plotting, a reasonably accurate curve has been constructed (Figure 1.51). Except for the points from the main blast line, the data does not show too great a variation.

The rositive duration curve is given in Figure 1.52. Again, as in the case of the arrival time data, all of the positive duration points are given in one curve. The spread of the data is small, and the curve can be assumed to be reasonably accurate. Correction factors for arrival times and positive duration were made in accordance with that given in Reference 14. The motcr of the self-recording gage requires 400 msec to establish its rated rpm. Correction factors of 65 msec for the $3-\mathrm{rpm}$ motor and 150 msec for the $10-\mathrm{rpm}$ motor were added to the arrival times obtained from the record. A nonlinear correction factor was applied to both time of arrival and positive duration where the true time of arrival was less than 400 msec .
1.3.2 Free-Field Blast Data For Various Projects. The free-field blast measurements taken for each project are given in Tables 1.9 through 1.16. A better proportion of pressuretime records was obtained from these gages than from the gages along the main blast line (Table 1.7).

Good pressure-time records were obtained for Project 1.7. The greatest variation in the maximum overpressure at equivalent ground ranges occurs at the 1,360 -foot station: one gage records a pressure of 40 psi and the other 60 psi . However, an inspection of the waveforms (Figure 1.23) indicates that the selection of the 40 -psi point resulted from the fact that the peak of that record was cut off. At the two closer stations, 760 feet and 1,040 feet, the

maximum overpressures are in better agreement (Figure 1.22).
Table 1.10 indicates that no pressures were recorded for three of the gages. These gages were inside the structures, and it can be assumed that no excess pressure above ambient entered into the structure. A small pressure was detected in the plenum chamber of Structure F3.3-9019.2, as indicated in Table i.11. The purpose and scope of these measurements are given in the reports prepared by the projects for which these measurements were taken.

The instrumentation provided Projects 4.3/33.2 and 6.1 resulted in separate blast lines. The blast line for Project 4.3/33.2 consisted of $P_{t}$-gages and q-gages extending from 2,030 feet out to 6,120 feet from ground zero. This line was parallel to the main blast line and approximately 300 feet south. The results from these gages are given in Tables !. 14 and 1.15.

The blast line for Project 6.1 was in a northeasterly direction from ground zero, and it consisted of $P_{t}$-gages only. The range for location of these gages extended from 1,250 feet out to 5,320 feet (Table 1.16). Except for Project 6.1, the free-field blast data taken for the other prcjects has been plotted in Figures 1.18 and 1.28. A separate overpressure versus distance curve was drawn for the values obtained for Project 6.1, and it is shown in Figure 1.53. The empirical data cf Figure 1.53 is yezy similar to that shown in Figure 1.18.
1.4 DISCUSSION

Although the data on overpressure versus time is not as complete as would be desired, a representative amount has been obtained, which leads to a better understanding of the magnitudes and waveforms of the pressure within a precursor zone. Certain characteristics of the pressure waveform, such as the inconsistencies in the magnitudes of maximum everpressure at equivalent ranges, although not resolved, can be attributed to several factors. These factors can be (1) acceleration effects due to ground motion, and (2) blast per se, (3) overshoot of the pressure capsule, and (4) turbulent flow created by the thermal layer ahead of the blast wave, and (5) presence of dust behind the blast wave. The latter, of course, is not experimenially verified. Inspection of the pressure-time waveforms points out the difficulty in the arbitrary selection of the maximum overpressure value, which, in effect, also contributes to the inconsistencies.

The sharp rise of the pressure front at the first two stations, 350 and 450 feet on the main blast line, implies that the precursor did not form as yet. At the next station, 650 feet, the pressure waveform is distorted; hence, the precursor was formed in the region between 450 and 650 feet. Typical precursor waveform patterns are observed as far as the $\mathbf{4 , 0 0 0}$-foot station. At the last station, 5,000 feet, the pressure wave has a sharp rise and is a classicaltype shock. The precursor zone, therefore, extended along the main blast line from approximately 450 feet out to 5,000 feet. On the other hand, the pressure waveiorms obtained for Project 4.3,'33.2 indicate a clearing up of the precursor at 3,939 feet, which is a classical-type shock wave. This anomaly is not readily explainable. The pressure magnitude at 4,000 feet along the main blast line is the same as the value at 3,939 feet of Project 4.3/33.2. Also, the positive duration and the time of arrival for both these stations are of the proper magnitude at least within the experimental error. It is highly unlikely that the difference in ground range between these two stations should indicate large variation in prcssure, time of arrival, or positive duration. The above anomaly leads to the belief that the shock was asymmetrical but not large enough to be observed except for the particular case above.

The relative merits of the outlined procedure given in Section 1.1.4 for determining the freestream dynamic pressure are difficult to evaluate. The difficulty lies principally in the degree of the dust influence on the pressure measurements. Lack of additional instrumentation precluded separating the dust portion of flow from air. If the preliminary estimates of the dust registry coefficient $n$ for the BRL $q$-gage can be believed, then the free-stream dynamic pressures as given are reasonably accurate. Further difficulty was encountered in the data reduction process. This arose when use was made of the ground surface pressure measurement for subtraction from the total head pressure to obtain the pressure difference. In some cases, the ground surface pressure time did not correspond to the total head pressure time. It was
then necessary, generally, to adjust the times of the total head pressure to correspond to the ground surface pressure time. Although there exists some question as to the validity of the calculated free-stream dynamic pressure, using the outiined procedure, the measurements taken during Shot Priscilla do increase the confidence for understanding the flow phenomena associated with precursor formation and also of being able to estimate the magnitudes of flow for dust and air combinations.
1.4.1 Predicted Maximum Ove:pressure Versus Distance Compared to the Measured Curve. The BRi- and AFSWP-predicted curves for overpressure versus distance as compared to the measured curves are shown in Figure 1.54. The measured curve compares favorably with the predicted curves. At the higher pressurc values, over 300 pss , the measured curve is about the same as the AFSWP-predicted curve. The BRL-predicted curve overestimated the values. Below 300 psi , the largest variation in the measured curve is approximately 20 percent as compared to the predicted curves.

This favorable comparison between predicted maximum overpressure and measured maximum overpressure indicates that the state of knowledge concerning maximum overpressure is good, especially since Shot Priscilla was the first event in which overpressures above 200 psi were measured.
1.4.2 Predicted Curve of Peak Dynamic Pressure Versus Distance Compared to Measured Curve. The comparison of the predicted dynamic with measured dynamic pressure is shown in Figure 1.55. In this figure, the predicted values are given for the ideal or free-stream dynamic pressure and for the pressure difference between total and static pressure as measured from previous operations. Comparison of the as-measured difference with the predicted curve gives a less favorable agreement than was obtained with the orerpressure-distance curves, except in the lower pressure region. The maximum difference between measured and pradirted is in the center portion of the curve. The construction of the predicted curve was based on a limited number of data points, ard use was made mainly of the SRI data. In retrospect however, an examination of the data points shown in Figure 1.3 indicates lower values of pressure from the BRL q-gage in the center portion of the curve. Since SRI data of past tests is above the BRL data, then the dust registry coefficient of the SRI gage may be higher than for the BRL gage, leading to higher values of pressure difference recorded. The BRL data of past tests is closer to the measured values of Shot Priscilla.

The corrected dynamic pressure or the calculated dynamic pressure versus distance approaches closer to the ideal curve than the uncorrected carve, but it is still about $\mathbf{5 0}$ to 100 percent higher. The extent of the dust influence, for Mach number calculations that are used for determining the corrected dynamic pressures as mentioned preriously, is unknown. Any discussion concerning the calculated dynamic pressures at this time would be premature. To discuss the validity of the applied corrections, additional information is needed principally in the dust registry coefficient of the gages used for measuring the flow phenomena.

### 1.4.3 Comparison of Predicted Values of Time of Arrival and Positire Duration with

 Measured Values. Shown in Figure 1.56 is a comparison of the predicted time of arriral and positive duration curves with the measured curves of Shot Priscilla. Except at the cioser ground distances, the comparison is reasonably good for both blast wave parameters. At closer ground distances the measured positive duration is shorter. and the time of arriral is longer than the predicted values. The predicted carve for positive duration. based on TM 23-200 (Reference ī). gives a better comparison with the measured curve that, the BRL-predicted curve. The BRL curre was drawn higher in the range of 1,500 to $\mathbf{5 , 0 0 0}$ fect because of thermal action predicted from Teapot 12 and Upshot-Knothole 10 data.Again as in the case of the overpressure-distance curce, the fair correlation betzeen the measured and predicted curves for time of arrival and positive daration indicates that the knowledge about basic blast phenomena is reasonably good.

### 1.5 CONCLUSIONS

Shot Priscilla was the first test during which emphasis was placed on results from the highpressure region. The overall free-field overpressure measurements yielded good pressuretime information, but the data on dynamic pressure is to some extent lin.ited. Loss of some of the dynamic pressure data resulted from high-acceleration effects.

An attempt had been made to apply appropriate correction factors to the q-sage data to obtain free-stream dynamic pressures. The validity of applying these correction iactors has not been ascertained. Lack of additiomal instrumentation precluded the separation of dust flow from airflow; therefore, the extent of dust influence on the measured data is unknown.

A good romparison was obtained between the predicted blast wave parameters and the measured values, except for dynamic pressure. The predicted dynamic pressures generally overestimated the measured values. However, a good correlation exists between the BRL data of the as-measured differential pressure from past shots with the values measured during Shot Priscilla. The predicted curve was based on past data obtained by SRI using an SC q-gage. The higher values of differential pressure from past tests by SC q-gage implies that the dust registry coefficient $n$ is higher for this gage than for the ERL q-gage. The calculated or corrected dynamic pressures are greater than the ideal by abou: 50 to 100 percent.

Upon final evaluation the new BRL q-gage does not agree too well with the old BRL q-zage.
The results during Shot Priscilla indicate that the state of knowledge concerning overpressure, time of arrival, and positive duration as a function of distance is good. The information obtained also leads to a better understanding of the flow phenomera associated with precursor formation and dust-laden sbock wares.

### 1.6 RECOMMENDATION

From results of past operations and this test, it is shown that the self-recording gages will gield reasomably accurate iniormation. Two of the shortcomings of these type gages are (1) the effects of acceleration, which can completely distort tie wareform and in some cases result in a complete loss of record, and (2) the lack of timing techaique to give more accurate time parameter measurements. A derelopment program is recommended to eliminate the abore two shortcomings of the gage.


TABLE 1.1 ASSUMED PARAMETERS AND SCALING FACTORS, SHOT PRISCILLA

| Location | Frenchman Flat | A-scaled factors |  |
| :--- | :---: | :--- | :--- |
| Elevation of ground zero, ft | 3,078 |  |  |
| Height of burst, ft | 700 | $\mathrm{~S}_{\mathrm{p}}{ }^{*}$ | $=1.150$ |
| Atmospheric pressure, ground zero, mb | 905 | $\mathrm{~S}_{\mathrm{d}}$ | $=0.2791$ |
| Atmospheric pressure, height of burst, mb | 881 | $\mathrm{~S}_{\mathrm{t}}$ | $=0.2765$ |
| Atmospheric temperature, ground zero, C | 17 | $\mathrm{~S}_{\mathrm{p}} \cdot \mathrm{S}_{\mathrm{t}}$ | $=0.3179$ |
| Atmospheric temperature, height of burst, C | 14.5 | $1 / \mathrm{S}_{\mathrm{p}}$ | $=0.8697$ |
| A-scaled height of burst, ft | 209 | $1 / \mathrm{S}_{\mathrm{d}}$ | $=3.583$ |
|  |  | $1 / \mathrm{S}_{\mathrm{t}}$ | $=3.617$ |
|  |  | $1 /\left(\mathrm{S}_{\mathrm{p}} . \mathrm{S}_{\mathrm{t}}\right.$ | $=3.146$ |

Pressure: $S_{p}=14.7 / P_{o}$
Distance: $\quad S_{d}=\left(P_{0} / 14.7\right)^{1 / 5}(1 / \mathrm{W})^{1 / 3}$

Time: $S_{t}=\left(T_{0}+273 / 288\right)^{1 / 2}\left(P_{0} / 14.7\right)^{1 / 3}(1 / W)^{1 / 3}$

TABLE 1.2 COMPARISON SHOTS AND SHOT PARAMETERS

| Shot | Yield | A-Scaled HOB | Shot Name |
| :---: | :---: | :---: | :---: |
| Greenhouse 2 | 46.7 | 82.9 | Easy |
| Tumbler 4 | 19.2 | 363.1 | - |
| Ivy 2 | 540.0 | 180.0 | King |
| Upshot-Knothule 1 | 16.5 | 112.5 | Annie |
| Upshot-Knothole 10 | $15.2 \pm 0.5$ | 203.4 | Grable |
| Upshot-Knothole 11 | $60 \pm 2$ | 316.8 | Climax |
| Teapot 2 | 2.39 | 213.4 | Moth |
| Teapot 3 | $6.8 \pm 0.2$ | 150.0 | Tesla |
| Teapot 4 | $43.2 \pm 2$ | 135.5 | Turk |
| Teapot 5 | $3.6 \pm 0.1$ | 186.1 | Hornet |
| Teapot 6 | $8.1 \pm 0.2$ | 240.1 | Bee |
| T eapot 8 | $14.2 \pm 0.7$ | 195.1 | Apple I |
| Teapot 10 | 3.2 | 13,195.7 | HA |
| Teapot 12 | 22.6 | 133.7 | Met |
| Teapot 13 | $28 \pm 1.5$ | 154.7 | Apple II |
| Teapot 11 | $28 \pm 1.5$ | 155.0 | Zucchini |
| Plumbbob 2 | $0.138 \pm 0.006$ | 561.8 | Franklin |
| Plumbbob 4 | $10.3 \pm 0.5$ | 218.3 | Wilson |
| Plumbbob 5 | $36.6 \pm 1$ | 201.7 | Priscilla |
| Plumbbob 6 | $71 \pm 2$ | 333.3 | Hood |
| Plumbbob 8 | $1.73 \pm 0.1$ | 11,708.0 | John |
| Plumbbob 9 | $10.3 \pm 0.5$ | 217.0 | Kepler |
| Plumbbob 10 | $9.7 \pm 0.5$ | 222.4 | Owens |
| Plumbbob 11 | $19 \pm 1$ | - | Stokes |
| Plumbbob 12 | $16.5 \pm 1.0$ | 184.5 | Shasta |
| Plumbbob 15 | $44+1$ | 187.6 | Smoky |
| Plambbob 16 | $11.1 \pm 1$ | 210.5 | Gallileo |
| Plumbbob 22 | $18.5 \pm 0.9$ | - | Whitney |
| Plumbbob 23 | $11.5 \pm 0.5$ | 633.0 | Charleston |
| Plumbbob 24 | $8.0 \pm 0.04$ | 237.5 | Morgan |

 DIRTY BLAST WAVES

```
\(q=\) dynamic air pressure \(=1 / 2 \rho u^{2}\)
\(q_{c}=\left(P_{p}-P_{s}\right)\)
\(M=u / c=\) local free-stream Mach number of flow behind blast front \(-\frac{\text { particle velocity }}{\text { sound velocity }}\)
\(P_{t}=\) free stream total pressure (absolute)
\(P_{p}=\) total head pitot pressure (absolute) \(\begin{aligned} & M<1 P_{p}=P_{t} \\ & M>1 P_{p} \neq P_{t}\end{aligned}\)
\(\mathrm{P}_{\mathrm{s}}=\) free stream static pressure (absolute)
\(\mathrm{P}_{\mathrm{O}}=\) ambient preshock static pressure (absolute)
\(\Delta P=\) free stream static overpressure \(=P_{S}-P_{0}\)
\(\Delta P_{p}=\) total head overpressure \(=P_{p}-P_{o}\)
\(\rho \quad=\) arr density (local)
\(u=\) particle speed of air (local)
c = speed of sound in air (local)
\(t=\) ratio of specific heats
Primes are used to denote uncorrected, "as read" gage values, thus
\(q_{c}{ }^{\prime}=\left(p_{p}-p_{s}\right)^{\prime}\)
Additional Symbols in Dirty Airblast Flows (*)
\(\mathbf{q}^{*}=\) dynamic air-plus-dust pressure in free stream \(=q+\phi_{d}\)
\(q_{c}^{*}=q_{c}+\phi_{d}\)
\(q_{c}^{* \prime}=\left(q_{c}+n \phi_{d}\right)^{\prime}\)
\(\phi_{d}=\) momentum flux of dust \(=\rho_{d} u_{d}{ }^{2}\)
\(\mathrm{n}=\) dust registry coefficient of gage, \(o \leq 2 . \cdot 1\)
\(\rho_{d}=\) mass of suspended dust per unit volume of mixture (local)
\(u_{d}=\) particle speed of dust (local)
\(\delta=\) specific gravity of dust particles \(=2.5\)
```



TABLE 1.6 Edget Berpromance

| Project | Pressure-Time Gages |  |  |  | Dynamic Pressuir Gages |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Total No. $\mathbf{P}_{\mathbf{t}} \text {-Gages }$ | Pressure-Time Records | Peak Pressure | No Record | Total No. of Gages | Press. $:$ re-Time Records |
| 1.1 | 22 | 14 | 6 | 2 | 15 | 12 |
| 1.7 | 9 | 7 | 2 | - |  |  |
| 3.1 | 6 | 6 | - | - |  |  |
| 3.2/3.3 | 6 | 4 | 2 | - |  |  |
| 3.4 | 6 | 3 | 3 | - | 4 | 3 |
| 4.3/33.2 | 7 | 5 | 1 | 1 | 7 | 7 |
| C. 1 | 15 | 11 | 3 | 1 |  |  |
| Totals | $\overline{71}$ | $\overline{50}$ | $\overline{17}$ | 4 | $\overline{26}$ | $\overline{22}$ |

TABLE 1.7 $\mathrm{P}_{\mathbf{t}}$-GAGE RESULTS, MAIN BLAST LINE

| Station | Ground Range | Maximum Overpressure | Arrival Time | Positive Duration |
| :---: | :---: | :---: | :---: | :---: |
|  | ft | psi | sec | sec |
| F1.1-9039.01A | 350 | No record |  |  |
| F1.1-9039.018 | 350 | 1030 | -- | -- |
| F1.1-9039-02A | 450 | 760 | -- | -- |
| F1.1-9039.023 | 450 | 750 | -* | 0.175 |
| F1.1-9039.03A | 650 | 480 | 0.364 | 0.95 |
| F1.1-9039.038 | 650 | 400 | 0.676 | 0,162 |
| F1.1-9040.01A | 850 | 225 | -- | 0.236 |
| F1.1-9040,018 | 850 | 206 | -- | -- |
| Fl.1-9040.02A | 1050 | 125 | -- | 0.233 |
| F1.1-9040.02B | 1050 | 138 | -- | 0.195 |
| F1.1-9041.00A | 2350 | 60.0 | -- | 0.343 |
| F1.1-9041.00B | 1350 | 62.0 | 0.512 | 0.280 |
| Fl.1-9042.01 | 1650 | 31.0 | - | 0.467 |
| P1.1-9042.02 | 2000 | 16.3 | -- | -- |
| F1.1-9042.05 | 2250 | 12.4 | 0.570 | 0.687 |
| F1.1-9042.06 | 2500 | 9.2 | 0.523 | 0.852 |
| F1.1-9042.07 | 3000 | 9.1 | -- | 0.727 |
| F1.1-9042.03 | 3500 | 9.9 | -- | - |
| F1.1-9042.08 | 4000 | 8.8 | 1.729 | 0.838 |
| F1.1-9042.04 | 4500 | 7.4 | -- | -- |
| F1.1-9043.01 | 5000 | 5.9 | -- | 0.916 |
| F1.19043.02 | 6000 | No r |  | -- |

TABLE 1.8 q-GAGE RESULTS, MAIN BLAST LINE, MAXIMUM VALUES

| Station | Ground Range | Total <br> Pressure | Static Overpressure | Pressure Difference $\left(\mathrm{P}_{\mathrm{p}}-\mathrm{P}_{\mathrm{o}}\right)^{* \prime}$ | Dynamic Pressure $\mathrm{q}^{*}$ | Mach Number |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | ft | psi | psi | psi | psi | (u/a) |
| F1.1-9040.01 | 850 | -- | -- | -- | -- | -- |
| F1.1-9040.02 | 1050 | 470.0 | 125.0 | 445.0 | 240.0 | 3.3 |
| F1.1-9041.00 | 1350 | 275.0 | 60.0 | 255.0 | 150.0 | 3.6 |
| F1.1-9041.00N ${ }^{\text {x }}$ | 1350 | -- | -- | -- | -- | -- |
| F1.1-9042.01 | 1650 | 143.5 | 31.0 | 150.0 | 80.0 | 2.3 |
| F1.1-9042.02 | 2000 | 58.5 | 23.0x | 44.0 | 32.0 | 1.3 |
| F1.1-9042.05N | 2250 | 48,0 | 12.4 | 36.0 | 27.0 | 1.4 |
| F1.j-9042.06 | 2500 | 47.0 | 9.2 | 38.0 | 25,0 | 1.3 |
| F1.1-9042.061ix | 2500 | 35.0 | 9.2 | 28.0 | 19.0 | 1.2 |
| F1.1-9042.07 | 3000 | 29.0 | 9.1 | 20.0 | 15.1 | 1.0 |
| F1.1-9042.07i. $x$ | 3000 | 26. 2 | 2.1 | 20.5 | 17.0 | 1.04 |
| F1.1-9042.03 | 3500 | 11.2 | $8.6 x$ | 3.4 | 2.8 | 0.45 |
| F1.1-9042.08 | 4000 | 10.0 | 9.0 | 1.3 | 2.3 | 0.29 |
| F1.1-9042.08N | 4000 | -- | -- | -- | -- | -- |
| F1.1-9042.04 | 4500 | 7.8 | 6.50 | 1.7 | 1.2 | 0.29 |

N , refers to new $q$-gage
$x$, values from q-gage

TABLE $1.9 \mathrm{P}_{\mathrm{t}}$-GAGE RESULTS, PROJECT 1.7

| Station | Ground <br> Range | Maximum <br> Overpressure | Arrival <br> Time | Positive <br> Duration |
| :--- | :---: | :---: | :---: | :---: |
| Fl.7-9031.01a | 760 | psi | sec | sec |
| Fl.7-9031.01b | 760 | 235 | 0.186 | -- |
| P1.7-9031.01c | 760 | 225 | 0.316 | 0.178 |
| Fl.7-9031.02a | 1040 | 112 | 0.186 | 0.126 |
| F1.7-9031.02b | 1040 | 115 | 0.250 | 0.307 |
| F1.7-9031.02c | 1040 | 110 | 0.241 | 0.253 |
| F1.7-9031.02d | 1040 | 105 | 0.306 | 0.256 |
| F1.7-9031.03a | 1360 | 600 | 0.210 | 0.285 |
| F1.7-9031.03b | 1360 | 400 | 0.305 | 0.404 |



| Station | Ground <br> Range | Maximum <br> Overpressure | Arrival <br> Time | Positive <br> Duration |
| :--- | :---: | :---: | :---: | :---: |
| F3.1-9014.01a | 860 | psi | sec | sec |
| F3.1-9014.01b | 860 | No pressure recorded | 0.260 |  |
| F3.1-9014.02a | 1040 | 118 | 0.082 | 0.087 |
| F3.1-9014.02b | 1040 | No pressure recorded | 0.254 |  |
| F3.1-9014.03a | 1360 | 56.1 | - |  |
| F3.1-9014.03b | 1360 | No pressure recorded | -- |  |

TABLE 1.11 Pt -GAGE RESULTS, PROJECTS 3.2 and $3.3 ~_{\text {-G }}$


TABLE 1.12 Pt -GAGE RESULTS, PROJECT 3.4


TABLE 1.13 q-GAGE RESULTS, PROJECT 3.4


TABLE $1.14 P_{t}$-GAGE RESULTS, PROJECT 4.3/33.2

| Station | Ground <br> Range | Maximum Overpressure | Arrival Time | Positive <br> Duration |
| :---: | :---: | :---: | :---: | :---: |
|  | ft | psi | sec | sec |
| F33.2-8015.01 | 2030 | 13.0 | $0.410^{\circ}$ | 0.010 |
| F3j.2-8015.02 | 2200 | 13.8 | 0.503 | 0.651 |
| F3j.2-6015.03 | 2750 | 9.0 | 0.903 | 0.757 |
| F35.2-6015.04 | 3450 | ¢. ${ }^{\circ}$ | 1.050 | 0.E2; |
| F35.2-6015.05 | 4770 | 4.3 | 2.300 | 0.920 |
| F35.2-8015.05 | 3320 | No S.ecord | -- | -- |
| F3j.2-8015.07 | 6120 | 4.9 | -- | -- |

TABLE 1.15 q-GAGE RESULTS, PROJECT 4.3/33.2, MAXIMUM VALUES

| Station | Ground Range | Total Pressure | Static Overpressure | Pressure Difference $\underline{1 D}_{\mu} \boldsymbol{P}_{s} i^{\prime=}$ | Dynamic <br> Pressure $\mathbf{q}^{*}$ | Mach number |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | ft | psi | psi | psi | psi | (u/a) |
| F35.2-8015.01 | 2030 | 61.0 | 12.0 | \%1.0 | 57.0 | 1.0 |
| Fj3.2-8015.02 | 2280 | 50.) | 13.8 | 41.0 | 28.0 | 2.3 |
| F33.2-8015.03 | 2750 | 23.7 | 9.0 | 14.5 | 11.5 | 0.89 |
| F3j.2-0015.04 | 3930 | 11.0 | 8.8 | 2.4 | 2.3 | 0.39 |
| F3j.2-8015.05 | 4770 | $6.0 \hat{0}$ | 6.5 | 0.9 | 0.6 | 0.23 |
| F53.2-8015.06 | -320 | 6.4 | $5.1 \times$ | 1.2 | 1.2 | 0.32 |
| F33.2-8025.07 | 6120 | 4.9 | $4.7 x$ | 0.3 | 0.35 | 0.10 ê |

$x$, obtained from q-gage as opposed to ground baffe gage.

TABLE $1.16 \mathrm{P}_{\mathrm{t}}$-GAGE RESULTS, PROJECT 6.1

| Station | Ground Range | Maximum Overpressure | Arrival Time | Positive Duration |
| :---: | :---: | :---: | :---: | :---: |
|  | ft | psi | sec | sec |
| F-6.1-1A | 2250 | 67.0 | -- | 0.422 |
| F-6.1-2A | 2370 | 57.0 | -- | 0.374 |
| F-6.1-3A | 1500 | 40.0 | -- | 0.510 |
| E-6.1-i4A | 1600 | 34.0 | -- | -- |
| F-6.1-5A | 1720 | 28.0 | 0.353 | 0.492 |
| F-5.1-6A | 1850 | 21.6 | 0.429 | 0.575 |
| F-5.1-7A | 1900 | 15.5 | 0.415 | 0.629 |
| F-6.1-8A | 2220 | 11.5 | 0.472 | 0.679 |
| F-ó.j-gA | 2290 | 11.0 | 0.535 | 0.751 |
| F-6.1-10A | 2520 | 10.5 | 0.609 | 0.830 |
| F-6.1-12A | 2730 | 8.5 | 0.772 | 0.808 |
| P-6.1-12A | 2870 | 10.0 | -- | -- |
| F-6.1-13A | 3250 | 8.0 | -- | 0.781 |
| F-0.1-i4 4 | 4530 | -- | -- | -- |
| F-6.1-15A | 5320 | 5.4 | $\cdots$ | -- |



Figure 1.1 Estimated peak overpressure versus distance, Shot Priscilla.


Figure 1.2 Data used in estimating overpressure-distance curve.


Figure 1.3 Estimated peak dynamic pressure versus distance.


Figure 1.4 Estimated time of arrival of initial disturbance.


Figure 1.5 Estimated time of arrival curve, computed curves, and experimental data.


Figure 1.6 Eistimated positiva-phase duriation.



Figure 1.8 Station layout, blast line, Shot Priscilla.



SWI- SAFETY SWITCH NORMALLY OPEN, CLOSED ',ANUALLY.
SWII- ROTATION LIMIT SWITCHES. TWO IN PARALLEL NORMALLY CLOSED.

Figure 1.12 Schematic diagram, photo-initiation circuit.


Figure 1.13 Turntable rotation versus elapsed time, 3-rpm motor.



Figure 1.15 Contractor-installed $q$-gage tower with new midbody for standard q-gage.

Figure 1.16 Contractor-installed towers for old and new model q-gages.


Figure 1.17 BRL-installed q-gage
mount, new model.

s.


Figure 1.18 Maximum overpressure versus distance, main blast line.


Figure 1.19 Overpressure-time histories, main blast line at distances of $350,450,650$, and 850 feet.

Figure 1.20 Overpressure-time histories, main blast line at distances of $1,050,1,350$, and 1,650 feet.



Figure 1.21 Overpressure-time histories, main blast line at distances of $2,250,2,500,3,000,4,000$ and 5,000 feet.


Figure 1.22 Overpressure-time histories, Project 1.7 at distances of 760 and 1,040 feet.

## 



Figure 1.23 Overpressure-time histories, Project 1.7 at distances of 1,040 and 1,360 feet.

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Figure 1．24 Overpressure－time histories，Projects 3.2 and 3.3 at distances of $970,1,040,1,150$ and 1,360 feet．





Figure 1.27 Overpressure-time histories, Project 6.1 at distances of $2,120,2,290,2,520,2,730$ and 3,250 feet.




Figure 1.28 Pressure difference $\left(P_{p}-P_{S}\right){ }^{* \prime}$ (maximum) versus distance.


Figure 1.29 Maximum dynamic pressure $q^{*}$ versus distance.




Figure 1.30 Pressure and Mach number versus time, main blast line at 1,050 feet.





Figure 1.31 Pressure and Mach number versus time, main blast line at 1,350 feet.



$0$





Figure 1.33 Pressure and Mach number versus time, main blast line at 2,000 feet.





Figure 1.34 Pressure and Mach number versus time, main blast line at 2,500 feet.







Figure 1.36 Pressure and Mach number versus time, main blast line at $\mathbf{3 , 0 0 0}$ feet. (1)


Figure 1.37 2ressure and Mach number versus time, main blast line at 3,000 feet.





Figure 1.42 Pressure and Mach number versus time, Project 3.4 at 4,200 feet.


Figure 1.43 Pressure and Mach number versus cime, Project 3.4 at 5,000 feet.






Figure 1.44 Pressure and Mach number versus time, Project 4.3/33.2 at 2,030 feet.
weracol






Figure 1.45 Pressure and Mach number versus time, Project 4.3/33.2 at 2,280 feet.


Figure 1.46 Pressure and Mach number versus time, Project 4.3/33.2 at 2,730 feet.



Figure 1.47 Pressure and Mach number versus time, Project 4.3/33.2 at 3,930 feet.






Figure 1.48 Pressure and Mach number versus time, Project 4.3/33.2 at 4,770 feet.





Figure 1.50 Pressure and Mach number versus time, Project 4.3/33.2 at 6,120 feet. Malian


Figure 1.51 Time of arrival versus distance.



Figure 1.52 Positive duration versus distance.


Figure 1.53 Maximum overpressure versus distance, Project 6.1.


Figure 1.ji Comparison of predicied maximum overpressure with measured values 252 function oi distance.


Figure 1.55 Comparison of predicted maximum dynamic pressure with measured values as a function of distance.


Figure 1.56 Comparison of predicted time of arrival and positive duration with measured values as a function of distance.
l
the separation between the two peaks shown in Figure 2.1 increases. In addition, the rise time for the main shock front also increases with distance. (The shape of the dynamic pressure wave need not follow that of the overpressure wave.) The wave is classified as Type A so long as the second peak is greater than the first peak pressure and so long as the time interval between arrival of the two maximums increases with distance from ground zero.

### 2.2 PREDICTION OF WAVEFORMS

The procedure for predicting wave type and pressure-time parameters is described below. The method amounts essentially to obtaining predicted values of type and of pertinent parameters of the pressure-time pulse from best-fit curves derived from data available from 12 precursor-forming blasts (Figures 2.3 through 2.12). These 12 shots were over several different desert-type surfaces; accordingly, the prediction technique is probably applicable only for shots over desert surfaces.
2.2.1 Prediction of Wave Type. Pressure-time data from the 12 shots has been (1) classified as to type according to the criteria established in Section 2.1 and (2) plotted on a scaled yield versus distance chart (Figure 2.3). With best-fit curves separating the four wave types completely, the chart shows that generally an accurate prediction of wave type can be made (for the distances and yield for which data is plotted).
2.2.2 Prediction of Type A Waveforms. Once the wave type is known, the prediction procedure is concerned with providing specific information about the wave, such as the ratio of first peak to second peak pressure. The procedure for predicting the specific data has been evolved only for the Type A waveform (a similar analysis has not yet been carried out for Types B and C waveforms). The Type A method is based on computation of the specific information shown in Figure 2.2 from the pressure-time records of the previously mentioned 12 shots. These computations have been plotted on various types of charts, and best-fit curves have been drawn (Figures 2.4 through 2.12). These curves, then, form the basis for prediction.

If the specifics of the wave shape are desired for a given yield, height of burst, and at a position where a particular peak overpressure exists, the first step in the prediction procedure is to establish the ground range corresponding to this pressure. This ground range can be obtained from charts such as Figures 2.4 or 2.5. Figure 2.4 gives the ground distances of specific overpressures for various heights of burst. Figure 2.5 gives the ground distance for a range of overpressures but for a given height of burst (around 200 feet).

Once the ground range of the desired overpressure has been established, the succeeding steps in waveform prediction involve using curves such as found in Figures 2.6 through 2.11 to obtain the sign'ficant wave parameters. There is considerable scatter of plotted points about the best-fit curves in many cases; this gives an estimate of the accuracy of predicted values. The best-fit curves can often be biased to give predictions for a particular shot of interest. For example, in the drawing of a curve all the data points for low shots might be ignored, thus making the curve more applicable to higher heights oi burst, which strongly indicates a height of burst dependency. This has been done in Figure 2.10 (TU 4, Upshot-Knothole ${ }^{1}$ and Teapot 12 were ignored). Finally, Figure 2.11 deserves some explanation, since the reliablity of the two drawn curves is not readily apparent. For an individual shot, such as Teapot 5, a trend can be seen for high values of the peak-pressure ratio to exist at great and small ground distances, with minimum values at intermediate distances. Almost all individual shots show this trend, and the curves have been drawn with this in mind.

The waveforms predicted for three pressure ranges for Shot Priscilla using the procedure described above are shown in Figure 2.13.

Tables 2.1 and 2.2 list the values and sources used to plot Figures 2.4 through 2.12. From the plotted figures it should be possible to predict the wave shape for any of the pressures in
the Type A classification occurring beyond a scaled ground range of 200 feet or 716 feet actual ground distance. In Figure 2.12, the slant range was usedxto plothrivait time, since some of the stations were close to ground zero.

### 2.3 PROCEDURE

To collect the most significant data for precursor waveform studies, the shots selected for participation were those with yields for which there was a paucity of data. The distances selected for location of gages were such that the complete precursor cycle would be covered. This is true particularly for surface pressure versus time. Because of the limited number of q-gages available, an extensive coverage of dynamic pressure versus time could not be made. On the other hand, it was expected that the selection of distances for placement of the $\boldsymbol{c}_{-}$-gages would yield sufficient data for the study of formation of the dynamic pressure waveform in the precursor zone cycle. Self-recording pressure-time and dynamic-pressure-time gages were used. Most gage stations were installed by BRL personnel.

Tabulations of the shots and the estimated yield on which participation was effected are given in Table 2.3. Included in the table are the distances from ground zero for pressure gage locations. Shots Priscilla and Smoky are included in Table 2.3, but the collection of data for precursor waveform studies was a secondary objective for these events. The data for Shot Smoky is reported in Reference 15. The field layouts of station locations for the shots, except Priscilla and Smoky, are shown in Figures 2.14 through 2.22.

### 2.4 RESULTS

To achieve the primary objective associated with this portion of the project, a total of 158 $\mathrm{P}_{\mathrm{t}}$ - and q-gages were deployed in nine shots. A small number of very-low-pressure (VLP) gages are included in the above total. The VLP data although not pertinent to this chapter is reported herein to extend the blast line information out to farther ranges from ground zero. Summaries of gage performance and shot data are given in Tables 2.4 and 2.5. Of the failures on the nine shots, some were due to the initiation system, and others were caused by highacceleration effects that broke the glass recording disks.

A large amount of data has been collected that will improve, upon final analysis, the present prediction technique for overpressure waveforms. The coverage of the dynamic pressure data was not as extensive as for the overpressure data. However, it is believed, that the data obtained, together with the data from past operations, will permit the development of a prediction technique for dynamic pressure waveforms.
2.4.1 Pressure-Time Histories. The plots of the pressure-time histories of all records are presented in Figures 2.23 through 2.58. The records are arranged by shot and by distance from ground zero, so the waveforms at discrete points along the blast times can be seen and compared to other shots. The maximum overpressure, arrival time, and positive phase duration associated with each $P_{t}$-gage record are listed in Table 2.6. The q-gage records are shown in Table 2.7. This table presents the maximum values for total and static pressure, pressure difference, and calculated dynamic pressure. Dynamic pressure records above 50 psi show excessive hash and high-frequency oscillation, whereas, in general, the static over-pressure-time records were good.

For stations where two $P_{t}$ gages were located, a reasonably good agreement was obtained for the maximum overpressure value and for the positive phase duration of the blast wave. The time of arrival records for some shots have considerable scatter; for other shots where no arrival time is given in the tables, the data was erratic and therefore meaningless. For the latter case, it was difficult to ascertain the start-up time because of a series of markings on the glass disk. The cause for the markings is unknown.

The records for the $q$-gages were reported only for those values that gave pressure-time histories. The reduction of the data was similar to that described in Chapter 1. It should be noted that no value for thic corrected dynamic pressure is presented for Shot Franklin. The selection of gage ranges and locations for Shot Franklin was based on a yield higher than actually realized. Hence, the deflections on the glass disks of the pressure-time record were small and possibly subject to large error in reading. For the low value of pressure difference obtained, the correction for Mach compressibility would be negligible, and for this reason is the existence of a possible large error, the application of a Mach correction factor did no. seem to be justified.
2.4.2 Waveforms. Precursor waveforms were observed on six of the nine shots. Of the six shots, Shot Galileo and Kepler showed only the early stages of the precursor; the remaining three shots showed only the ideal Type D waveform. Between Shots Galileo and Kepler, the precursor waveform of Shot Kepler formed into an ideal waveform sooner than during Shot Galileo.

The nonprecursor shots were Franklin, Shasta, and Charleston. For Franklin, as mentioned previously, gage ranges and locations were selected based on a higher yield. If a gage had been located at a range closer to ground zero, it is possible that a precursor would be evident. On the cther hand, the height of burst for Shot Franklin is near the region where a precursor would not form, based on the criteria of Figure 2.17 of Reference 7. Simiiarly, the burst height of Shot Charleston is in the same category as Shot Franklin. Hence, the lack of precursor for these two events could be predicted. Based on the same criteria of Refercnce 7, a precursor should have been evident during Shot Shasta.

The classification of waveform for each event is presented in Table 2.6. All the precursorproducing shots provided waveform data in the desired region above 50 psi . Shot Hood also fulfills the requirement for more waveform data at scaled heights of burst greater than 300 feet and for larger yields than previously considered.
2.4.3 Maximum Pressure versus Distance. Curves showing the maximum pressure values as a function of distance are given in Figures 2.59 through 2.73. Along with the dynamic pressure $q^{*}$ plots, the values of the difference between total and static pressure are also plotted. The maximum overpressure versus distance curves generally confirm the absence or presence of precursor formation and also indicate the degree of precursor actior For Shots Wilson and Hood (Figures 2.61 and 2.63), a typical precursor depression of the curve is observed. The depression of the curve for Shot Hood extends out to farther ranges and to lower pressure values than the curve of Shot Wilson, indicating a stronger precursor action during Shot Hood. The curves for Shots Franklin and Kepler (Figures 2.59 and 2.67) show the steady decrease of pressure with distance, which is the usual pattern for nonprecursor conditions.

A trend is apparent in the dynamic pressure-distance curve in the presence of a precursor. Following a steady decrease of pressure with distance, the pressure decreases gradually, then drops steeply when the precursor no longer exists. This is shown in Figure 2.62 for Shot Wilson and Figure 2.66 for Shot Owens. The above characteristics are also found in the curve for Shot Hood (Figure 2.64) except for the later portion of the curve. The steep decrease of pressure with distance for Shot Hood is not as pronounced as it is for Wilson and Owens.
2.4.4 Arrival Time and Positive Duration. The arrival time and positive duration measurements are plotted in Figures 2.74 through 2.82. The curves drawn through the data points were fitted by eye. The arrival time curves are less accurate than the positive duration curves. The data is useful to add to the present state of knowledge about nuclear blast waves, particularly the data from Shot Hood (Figure 2.76) where little information is available for this height of burst.


### 2.5 DISCUSSION

2.5.1 Nonprecursor Shots. The shot that did not yield precursor waveforms (although expected) was Shasta. In the early stages of Shots Kepler and Galileo, precursor waveforms were obtained, but they very quickly formed into ideal waveforms. For Kepler and Galileo, the usual precursor cycle was not observed in the same pattern as for Shots Wilson or Hood. It is of interest to note that Kepler, Shasta, and Galileo were tower shots with the devices heavily shielded, and Wilson and Hood were balloon shots with no appreciable shielding. It is felt that these heavy shields minimized the factors for precursor formation.

The general layouts of shield configurations for Shots Kepler, Galileo, and Shasta are shown in Figures 2.83 through 2.85. The long axis of the shield for Kepler and Galileo was approximately parallel to the blast line. The growth of the fireball for the heavily shielded shots would reveal whether the shields affected the thermal output of the detonation, and any affect of the thermal output would, in turn, reflect in the formation of precursors. This premise is based on the fact that an experimental thermal pulse can be calculated from the fireball radii as a function of time.

Motion photography of the fireball growth indicates that, during Shots Kepler and Galileo, the radius along the blast line as a function of time is less than in the direction perpendicular to the blast line (Reference 16). For Shot Kepler, the difference in fireball radius, from photography, expressed in terms of the ratio of weapon yield gives:

$$
\begin{equation*}
\frac{W_{1}}{W_{2}}=\left[\frac{\phi_{1}}{\phi_{2}}\right]^{5}=1.22 \tag{2.1}
\end{equation*}
$$

Where: $W_{1}=$ weapon yield, kt, perpendicular to the blast line
$W_{2}=$ weapon yield, kt, along blast line
$\phi_{1}=\left(d_{1} / t^{2 / 5}\right)$
$\phi_{2}=\left(d_{2} / t^{2 / 5}\right)$
and $\mathrm{d}=$ radius of fireball, meters
$\mathrm{t}=\mathrm{time}, \mathrm{msec}$
The yield of weapon $W_{2}$ along the blast line is less by about 22 percent. The fireball growth as a function of time of Shot Newton, a nonshielded device and similar in yield to Shot Kepler, compares reasonably well with Shot Kepler. A small variation in radius will indicate a large variation in yield because of raising the value $d$ to the fifth power. The variation in yield by itself is not sufficient to show any affect on precursor formation.

The irradiation rate $\mathrm{dq} / \mathrm{dt}$ from the fireball is given by:

$$
\begin{equation*}
\frac{d q}{d t}=f \sigma T^{d} 4 \pi R^{2} \tag{2.2}
\end{equation*}
$$

Where: $f=$ fraction of radiation spectrum transmitted
$\sigma=$ Stefan-Boltzmann constant
$T$ = teinperature of fireball
$\mathbf{R}=$ radius of fireball
For Shot Kepler, then, assuming all values to be constant except $R$ : the radiation rate difference along the blast line is 8 percent less than it is perpendicular to the blast line. This small decrease of the thermal pulse along the blast line would not minimize the thermal action for precursor formation. Evidently, for Kepler, the fireball growth was approximately normal, and the effect of minor variation of the fireball on precursor formation, from the above discussion, does not appear to be justified. On the other hand, consideration was not given to two parameters for the irradiation rate from the fireball radius; one is the value $f$ and the other is the temperature T .

The shield consisted of large masses of steel, concrete, paraffin, and lead. This, upon decomposition by the intense heating, would result in a considerable amount of debris in the surrounding vicinity other than air. This debris would then decrease the amount of thermal energy transmitted and also decrease the temperature of the fireball. Evidence of this decrease is indicated to some degree by the measurements of thermal radiation phenomena by Project 5.3 (Reference 17). Reference 17 contains a comparison between the measured and calculated quantities of thermal radiation. The measured thermal energy output was 100 to 400 percent less than the calculated values for the shielded shots and about 30 percent less for the nonshielded shots. Furthermore, between the calculated and measured values of time to maximum temperature rise, the measured valıes were larger by greater percentages for the shielded shots than for the nonshielded shots. The decrease of thermal radiation and the increase of time for the maximum temperature rise would decrease the heating of the surface prior to shock arrival and, in turn, minimize the conditions for precursor formation. Qualitatively then, it can be inferred that the shielding of the device will minimize the thermal action for precursor formation. From a military standpoint the above discussion has no real significance, since it is unlikely that any weapon will be fired with a large mass of shielding about the device. Therefore, it can be implied that for the shot conditions of Kepler, Galileo, and Shasta, a precursor would be formed without the existence of the shield around the device.
2.5.2 Scaled Pressure-Distance Curves. The scaled maximum overpressure-distance curves are shown in Figures 2.86 through 2.94 . Points have been selected from these scaled curves (including precursor and nonprecursor shots) and plotted on height-of-burst curves versus ground range (Figure 2.95). Not many of the points lie near the isopressure contours given in this set of height-of-burst curves. For those shots at lower burst height and higher pressures, a reasonable fit is obtained. However, for the higher burst heights and lower pressures, a large difference exists. In view of the additional data obtained and the variations noted, the height-of-burst curve as presently constructed should be revised. The dynamic pressure versus distance curves have not been scaled or compared to any height-of-burst curves. In view of the data accumulated in this operation and the new technique used for data processing and analysis, new height-cf-burst curves versus ground range are needed for dynamic pressure as well.

Of interest is the comparison of the basic blast parameters between precursor and nonprecursor shots. Wilson, a precursor snot, was selected for comparison with Kepler, since the scaied neight of burst and weapon yield are similar for both events. Figure 2.96 compares the overpressure for Shots Kepier and Wilson. The difference in pressure at the close-in stations is not as great as at the farther ranges. For the precursor shot, the usual depression in the pressure-distance curve is observed, whereas there is a steady decrease of pressure with distance for the nonprecursor shot. Lack of data at closer stations for Shot Wilson precludes sny comparison of dynamic pressure as a function of distance (Figure 2.97). A later time of arrival is noted for Shot Kepler than for Shot Wilson (Figure 2.98). For strong precursor action, as during Shot Wilson, this would be expected. Also, another distinction between precursor and nonprecursor shots, is the longer duration for precursor shots (Figure 2.99); the positive duration is longer for Shot Wilson than for Shot Kepler.

The overpressure values obtained on the nonprecursor shots are compared with a predicted curve for the good surface height-of-burst curve of TM 23-200 (Reference 7). In general, a reasonable comparison is obtained particularly at the lower pressure values (Figure 2.100:-
2.5.3 Wave-Type Prediction Chart. Wave types from precursor-forming shots have been plotted on the BRL wave-type-prediction chart (Figure 2.101). It can be seen from the ploted data that there is general agreement with data from previous shots, except in the case of Si:ot Hood. The Hood data is particularly interesting, as it is at a scaled height of burst for which: data was needed, and because it is apparent that the current prediction chart will have to be medified for yields and heights of burst of the order of those of Shot Hood.

2.5.4 Waveform Prediction. It can be seen from Figure 2.101 that the Priscilia waveform data in the high-pressure region, Type A, falls in an area where little previous data was ayailable. It was, therefore, of interest to see whether the waveform prediction technique outlined in Section 2.2.2 would be valid for this region. Predictions were mace for Shot Priscilla for the two peak overpressures, 70 and 40 psi . Good overpressure-time records were obtained from Priscilla for these pressure levels. The actual waveforms have been drawn in Figure 2.102 and are compared to the predicted waveforms. The comparison shows gocd agreement throughout.

### 2.6 CONCLUSIONS

Overpressure wave data was collected from six precursor-forming shots (not including Priscilla and Smoky), extending the capability for predicting waveforms. In particular, waveform data in pressure regions above 50 psi was furnished by Shots Hood, Wilson, and Owens. Aiso, Shot Hood provided data for scaied heights of burst greater than 300 feet.

* Considerable dynamic-pressure waveform data was collected from the three precursorforming shots. This data was of good quality for prediction-development pirposes in overpres--are regions below 50 psi , but its quality in regions above 50 psi is questionable.
EThe current BRL wave-type-prediction method appears to be applicable to yields lower than $601+$ and to heights of burst ranging from 136 to 300 feet. At higher yields or scaled heights of burieka revision in the prediction chart is necessary.

The:current BRL waveform-prediction method for Type A appears to be valid for overpressure einges up to 70 ps , at least for heights of burst less than 300 feet, and yields less than 60 14

Twof four tower shots (not including Smoky) gave evidence of precursor waveforms that weréshothived. All balloon shots produced strong precursor zones except Shot Charleston. The two hiots, one tower and one balloon, were of such a nature that precursors were not expectedtoform: according to the chart given in Reference 7. It is felt that heavy shielding of the nuclemedevice decreaced the thermal action for precursor formation.

Comparionof measur ed overpressures with the isopressure contours on a height of burst versus groundirange indicate large variations at the lower pressure values.

### 2.7 RECOFMENDTION

The recomimendation made in Chapter 1 is applicable to this portion as well. Although menin nas been mbde for the need of new height-of-burst curves for overpressure and dynamic -ssure, sunh with is already under progress or has been completed at the time of writing on s report.

TABLE 2.2 TYPE A PRECURSOR DATA. A-SCALED

| Slant <br> Range | Horizontal Distance | ${ }^{1} 1$ | ${ }^{1}+$ | $\mathrm{tr}_{r}$ | $\mathrm{t}_{\mathrm{ps}}$ | $t_{\text {m }}$ | $\mathrm{P}^{1} \mathrm{p}_{1} / \mathrm{l}^{\prime} \mathrm{m}$ | $\mathrm{P}_{\mathrm{P}_{2}} / \mathrm{Pm}_{\mathrm{m}}$ | $\mathrm{P}_{\mathrm{p}_{2}} / \mathrm{P}_{\mathrm{p}_{1}}$ | $\mathrm{P}_{\mathrm{m}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| feet | feet | sec | sec | sec | sec | sec |  |  |  | psi |
| Dishot-Knothole Shot 10, Yield 14.9 kt *. Sealed Height of Burst 203 feet. $\mathrm{S}_{\mathrm{d}}$ 0.3885, $\mathrm{S}_{\mathrm{t}}-0.3839, \mathrm{~S}_{\mathrm{p}}=1.146$ |  |  |  |  |  |  |  |  |  |  |
| 323 | 26: | 0.052 | 0.073 | 0.002 | 0.009 | 0.020 | 0.17 | 0.14 | 0.80 | 159.0 |
| 360 | 297 | 0.062 | - | 0.006 | 0.011 | 0.015 | 0.21 | 0.40 | 0.95 | 128.3 |
| 393 | 336 | 0.06 C | 0.10i | 0.010 | 0.018 | 0.024 | 0.17 | 0.01 | 0.07 | 73.1 |
| 402 | 347 | 0.068 | 0.108 | 0.003 | 0.023 | 0.028 | 0.17 | 0.06 | 0.36 | 80.0 |
| 412 | 358 | 0.07 .1 | 0.105 | 0.010 | 0.023 | 0.027 | 0.20 | 0.19 | 0.95 | 81.9 |
| 498 | 454 | 0.102 | 0.123 | 0.011 | 0.040 | 0.092 | 0.48 | 0.21 | 0.45 | 43.1 |
| 585 | 551 | 0.135 | 0.175 | 0.007 | 0.069 | 0.077 | 0.072 | 0.29 | 0.41 | 16.4 |

Upshot-Knothole Shot 11. Yield 60.5 kt . Scaled Height of Burst 316 feet. $\mathrm{S}_{\mathrm{d}} \quad \mathbf{0 . 2 3 7 0}, \mathrm{S}_{\mathrm{t}} \quad 0.2340, \mathrm{~S}_{\mathrm{p}}=1.229$

| 455 | 363 | 0.111 | 0.136 | 0.001 | 0.023 | 0.034 | 0.25 | 0.08 | 0.33 | 48.7 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 572 | 477 | 0.145 | 0.160 | 0.002 | 0.041 | 0.057 | 0.35 | 0.25 | 0.73 | 29.8 |

Cpshot-Knothole Shot 1. Yield 16.2 kt . Sealed Height of Burst 112 feet, $\mathrm{S}_{\mathrm{d}}=0.3750 \quad \mathrm{~S}_{\mathrm{t}}=0.3670 . \mathrm{S}_{\mathrm{p}}=1.1704$

| 286 | 262 | 0.032 | - | 0.004 | 0.007 | 0.013 | 0.54 | 0.22 | 0.41 | 131.1 |
| ---: | ---: | ---: | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 356 | 335 | 0.044 | - | 0.005 | 0.018 | 0.031 | 0.59 | 0. | 0.53 | 56.6 |
| 410 | 394 | 0.056 | 0.146 | 0.009 | 0.029 | 0.049 | 0.70 | 0.42 | 0.60 | 38.6 |
| 446 | 431 | 0.066 | 0.141 | 0.005 | 0.044 | 0.060 | 0.91 | 0.26 | 0.29 | 22.8 |

Shot TE 4. Yield 19.6 kt . Scaicd Height of Burst 363 feet. $\mathrm{S}_{\mathrm{d}}=0.349 . \mathrm{S}_{\mathrm{t}}=0.346 . \mathrm{S}_{\mathrm{p}}=1.199$

| 421 | 212 | 0.098 | 0.104 | 0.001 | 0.006 | 0.012 | 0.33 | 0.30 | 0.91 | 99.2 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 497 | 340 | 0.127 | 0.143 | 0.005 | 0.018 | 0.026 | 0.24 | 0.23 | 0.95 | 52.3 |
| 594 | 469 | 0.168 | 0.186 | 0.007 | 0.042 | 0.056 | 0.36 | 0.37 | 1.62 | 29.4 |

Shot TP 2. Yield 2.4 kt . Scaled Height of Burst 213 feet. $\mathrm{S}_{\mathrm{d}}=0.7112 . \mathrm{S}_{\mathrm{t}}=0.6812 . \mathrm{S}_{\mathrm{p}}=1.163$

| $\mathbf{5 2 1}$ | $\mathbf{4 5 5}$ | 0.100 | 0.129 | 0.011 | 0.025 | $0.03:$ | 0.45 | 0.14 | 0.21 | 56.3 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $\mathbf{c o 3}$ | 564 | 0.111 | 0.150 | 0.007 | 0.043 | 0.056 | 0.75 | 0.43 | 0.55 | 26.7 |

Shot TP 4 . Yield 43 kt . Scaicd ifcight of Burst 136 feet. $\mathrm{d}=0.26 y 6 . \mathrm{S}_{\mathrm{t}}=0.2630 . \mathrm{S}_{\mathrm{p}}=1.136$

| 495 | 480 | - | 0.169 | 0.007 | 0.072 | 0.077 | 0.57 | 0.33 | 0.53 | 32.0 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |

Shot TP 5. Yield 3.6 kt . Scaled Height of Burst 196 feet. $S_{d}=0.6203 . S_{t}=0.6064 . S_{p}=1.161$

| 433 | 391 | 0.065 | 0.090 | 0.002 | 0.021 | 0.023 | 0.35 | 0.03 | 0.23 | 88.2 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 479 | 441 | 0.073 | 0.097 | 0.004 | 0.035 | 6.041 | 0.60 | 0.11 | 0.15 | 51.1 |
| 354 | 522 | 0.953 | 0.152 | 0.008 | 0.045 | 0.059 | 0.30 | 0.13 | 0.16 | 34.8 |
| 613 | 584 | 0.099 | 0.170 | 0.003 | 0.072 | 0.076 | 0.97 | 0.32 | 0.33 | 19.7 |

Shot TP 6. Yield 7.76 kt . Scaled Height of Burst 240 fect. $\mathrm{S}_{\mathrm{d}}=0.4593 . \mathrm{S}_{\mathrm{t}}=0.4679 . \mathrm{S}_{\mathrm{p}}=1.163$

| 572 | 519 | 0.124 | 0.125 | 0.010 | 0.050 | 0.065 | 0.55 | 0.42 | $0.5 S$ | 22.3 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |



| $4 i t$ | 402 | 0.075 | 0.103 | 0.005 | 0.031 | 0.033 | $0.5 i$ | 0.31 | 0.57 | 35.9 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |



| $\because 91$ | 257 | 0.036 | - | 0.007 | 0.005 | 1.010 | 0.49 | 0.29 | 0.59 | 155.6 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 369 | 313 | 0.051 | 0.123 | 0.002 | 0.016 | 0.029 | 0.56 | 0.12 | 0.21 | 77.6 |
| 532 | 517 | 0.021 | 0.189 | 0.003 | 0.062 | 0.055 | 0.66 | 0.17 | 0.26 | 33.5 |
| 532 | 517 | 0.07\% | 0.144 | 0.007 | 0.073 | 0.079 | 0.55 | 0.33 | 0.35 | 30.6 |

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| 4.35 | 390 | 0.07t | 0.120 | 0.004 | $0.0 \geq 6$ | 0.036 | 0.34 | 0.06 | 0.15 | S0.5 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 49! | 40 | 0.090 | 0.135 | 0.00; | 0.035 | 0.043 | 0.76 | 0.07 | 0.15 | 57.5 |
| 325 | ;90 | 0.105 | 0.155 | 0.095 | 0.08i | 0.060 | 0.55 | 0.15 | 0.36 | 10.? |
| 602 | 5.0 | 0.152 | 0.154 | 0.013 | 0.072 | 0.090 | 0.90 | 0.75 | 0.56 | 23.0 |

[^0]
## 

Pable: as gadie tocations for brecursor waverolm stumes Gage locations are givon un distmatas from ground zoro in foot.

| Shot | Fruaklin | Wilson | Prinella | Howl | Kupler | Owens | Shasta | Smoky | Gallioo | Charloston | Morgan |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Eistimated steld. it | 0.138 0.0001 | 10.315x | 10,641 | 74.1 $60 \%$ | 10.316.5 | 0.0 4.5 | 16.815\% | 43.71:5\% | 11.311 | $11.5 \pm 5 \%$ | 7.8 $\pm 5 \%$ |
| Helfith of burst. it | 30014 | 60013 | 70013 | 1,60013 | 600'1 | 50013 | $5001{ }^{\circ}$ | 700'1 | 500 T | 1,50013 | 50013 |
|  | 100 | (i00 ¢ | 350 | 1,000 | 550 | 500 | 1.000 | 8.40 | 600 | 1,100 | - |
|  | 700 | 1.000 c | 460 | 1,011 | 000 | 10009 | 1.500 | 800 | 1,050 | 1,300 | 600 |
|  | 8009 | 1,100 | 050 | 1.700 C | 1,050 | 800 | 2,000 | 1,005 | 1,200 9 | 1,500 | 1,000 |
|  | 000 al | 1.2009 | 8509 | 2.2009 | 1,200 ${ }_{\text {c }}$ | 1,000 ¢ | 2.500 | 1,170 | 1,100 | 2,300 | 1,250 |
|  | 1.000 | 1.420 | 1.0509 | 2.100 | 1.400 | 1.150 | 3,000 | 1,5009 | 1,700 9 | 3.000 | 2,200 |
|  | 1.2009 | 1.011 | 1.1509 | 3,0009 | 1.700 q | 1.120 | 18.000 | 2.018 | 1,050 | 5,500 | 3.000 |
|  | 1,000 11 | 1,076 | 1.050 q | 4.0004 | 1.050 | 1,511q |  | 2,5809 | 2,100q9 | 8,000 | 5,000 |
|  | 1,800 | 1.000 | 2,000 1 | 4.600 | 2.1009 | 1.575 |  | 2,0439 | 2,200 | 16,000 | 6,000 |
|  |  | 1,7004 | 2.2009 | 0.00011 | 2,200 | 1,060 |  | 3.4009 q | 2,350 | 21.120 | 8.000 |
|  |  | 1,800 | 2.6009 | 0.600 | 2.350 | 1.700 q |  | 3,8759 | 2,600 | 25.872 | 15,000 |
|  |  | 1.900 g | 3.000 q | 6,000 | $2,500 \mathrm{q}$ | 1,800 |  | 4,165 | 2,600 | 51.216 | 21,120 |
|  |  | 2,100 | 3.0009 | 0.6009 | 2,000 | 3,000 |  | 4,320 | 2,7659 | 73.000 | 25,872 |
|  |  | 2.2009 | 4.0009 | 7.000 | 2,800 | 0,000 |  | 5.680 C | 4.700 |  | 35,000 |
|  |  | 2.100 | 1.600 | 8.800 | 22.016 |  |  | 0,000 |  |  |  |
|  |  | 2,800 | 6,000 |  |  |  |  | 8,000 |  |  |  |
|  |  | 3.000 | 13.000 |  |  |  |  |  |  |  |  |

Tablet: a.l Ginge: perpomatance:

| Shot | Total Numiver $r_{1}$-Chagen | $\begin{aligned} & \text { Pressure-TIne } \\ & \text { Menord } \\ & \text { hatere } \end{aligned}$ | Guges <br> Peak Pressute Only | No Rucord | Dynamice Pr $\substack{\text { Total Number } \\ \text { q-Gagos }}$ | $\qquad$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Prankiln | 11 | 0 | $:$ | $\cdots$ | 4 | 4 |
| Wilnote | 18 | 17 | 1 | - | 8 | 5 |
| Higend | 18 | 1.1 | 4 | - | 1 | 0 |
| - Worits | 16 | 12 | 4 | - | 1 | 4 |
| Kepler | 16 | 1.1 | 1 | - | 5 | 3 |
| stemia | 6 | 5 | 1 | - | - | - |
| Caliluo | $1: 1$ | 12 | 1 | - | 4 | 4 |
| Citailontern | 18 | 10 | 6 | 2 | - | - |
| Musgall | 12 | 1 | 3 | - | - | - |
| 'foldt. | 127 | 102 | 23 | 2 | 31 | 45 |


$\square$


| Station |  |  | $\begin{aligned} & \text { Arrivel } \\ & \text { yive } \end{aligned}$ | $\begin{aligned} & \text { Pometive } \\ & \text { Duration } \end{aligned}$ | Type iarefors |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | sret | pel | see | sec |  |
| geot 12millin |  |  |  |  |  |
| 3-1.1-9004.00 | 400 | 25,0 | 0.585 | 0.066 | D |
| 3-1.1-500 +.00 | 400 | 26.0 |  | - |  |
| 3-1.1-9004.01 | 700 | 12.8 | 0.269 | 0.098 | D |
| 3-1.1-900t.01 | 200 | 12.2 |  | 0.095 | D |
| 3-1.1-9003.01 | 800 | 7-7 | 0.328 | 0.092 | 7 |
| 3-1.1-9003.01 | 800 | 8.6 | 0.359 | 0.165 | 1 |
| 3-1.2-5003.02 | 900 | $7-2$ | 0.202 | 0.000 | D |
| 3-3 -1-500t.ce | 1200 | 6.4 | 0.202 | 0.153 | D |
| 3-2.1-9003.03 | 1200 | 4.8 | 0.368 | 0.153 | D |
| 3-1.1-5003.03 | 1500 | 3.5 | 0.568 | 0.122 | D |
| 3-1.2-900 -03 | 1800 | 2.5 | - | - | - |
| sut 4sizuen |  |  |  |  |  |
| 9-1.1-9921.01 | 600 | 110.0 | 0.405 | 0.144 | A |
| 9-1.1-g0en 01 | 600 | 140.0 | 0.036 | 0.170 | 4 |
| 9-1.1-902.02 | 3000 | 38.0 | 0.119 | 0.276 | 4 |
| 9-1.1-9024.02 | 1000 | 12.0 | 0.294 | 0.308 | 4 |
| 9-1.1-90e2.01 | 1120 | 31.0 | - | - | - |
| 9-1.1-902n.03 | 120 | 240 | 0.137 | 0.392 | 4 |
| 9-1.1-902.031 | 170 | 12.0 | 0.200 | 0.329 | 3 |
| 9-1.1-902.02 | 1512 | 10.6 | 0.200 | 0.196 | - |
| 9-1.1-50es an | 1575 | 12.6 | - | 0.176 | B |
| 9-1-1-s0e2.06: | 1620 | 90 | 0.430 | 0. 136 | B |
| 9-1.1-902.01 | 1700 | 9.3 | 0.48 | 0.420 | B |
| 9-1.1-sper.03 | 3190 | 9.0 | 0.499 | 0.49t | 3 |
| 9-1.1-9021.05 | 1980 | 9.7 | 0.659 | 0.102 | c |
| 9-1.1-902.06 | 2100 | 30.1 | 0.67 | 0.402 | c |
| 9-1.1-9024.01 | 2000 | 9.3 | 0.076 | ‥22 | c |
| 9-1.1-90er.05 | 2100 | 7.8 | 1.003 | 0.533 | D |
| 9-1.1-9021.06 | 2600 | 6.6 | 1.063 | 0.59 | D |
| 9-1.1-9021.08 | 300 | 6.2 | 1.437 | c. 613 | D |
| Stor Eoxd |  |  |  |  |  |
| 9-1.1-9001.00 | 1000 | 68.0 | - | 0.406 | 4 |
| 9-1.1-9001.00 | 1000 | 92.0 | - | 0.465 | 4 |
| 9-1.1-9022.02 | 1511 | \$2.0 | 0.394 | 2.545 | 4 |
| 9-1.1-9001.04 | 1700 | 4,5 | 0.721 | 0.709 | 4 |
| 9-1.1-9001.07 | 2200 | 18.7 | 0.543 | 0.838 | A |
| 9-1.1-9002.05 | 2300 | 17.0 | 0.65 | 0.603 | A |
| 9-1.1-9001.06 | 3000 | 8.5 | 0.853 | 1.9is.3 | 3 |
| 9-1.1-9021.09 | +000 | 6.7 | - | - | - |
| 9-1.1-900:.09 | 1000 | 6.6 | - | - | - |
| 9-1.1-9022.07 | -590 | 6.8 | 1.76? | 1.153 | $c$ |
| 9-1.1-9001.10 | 5000 | 5.2 | 2.260 | 2.173 | c |
| 9-1.1-9002.06 | 5500 | 6.0 | 2.670 | 1.185 | D |
| 9-1.1-9002.09 | 8000 | 5.4 | 2.838 | 2.229 | D |
| 9-1.1-9002.09 | 5000 | 5.5 | 2.04 | :-272 | D |
| 9-1.1-902:.11 | 5600 | *. 7 | 3.352 | 1.35d | 5 |
| 9-1.1-9002.10 | 7000 | 4.1 | - | - | - |
| 9-1.1-908.12 | 0000 | 3.5 | \$.118 | 1.26 | D |
| 9-1.1-9002.11 | 0000 | 5.3 | - | - | - |
| Stot Ieple |  |  |  |  |  |
| L-E-1-9020.00 | 350 | 195.0 | - | 0.ili | 4 |
| 4-1.1-3001.00 | 900 | 35.0 | - | 0.185 | 4 |
| 1-1.1-9021.01 | 1050 | 32.0 | 0.198 | 0.220 | 3 |
| \%-1.1-9021.01 | 1050 | +5.0 | 0.140 | 0.235 | D |
| i-1.1-9002.0i | 1200 | 30.0 | 0.256 | 0.253 | D |
| 1-1-1-9021.00 | 1400 | 25.0 | 0.216 | $0.3+5$ | D |
| 4-1.1-9022.02 | 1700 | 20.0 | 0.76 | 0.381 | 0 |
| +1.1-9021.03 | 1550 | 1\%. 5 | 0.575 | 0.41 | 0 |
| t-1.1-9022.03 | 2100 | 15.0 | 0.735 | 9. 153 | $p$ |
| 4-1.1-9001.0t | 2200 | 12.3 | 0.70 | 0.630 | D |
| +1.1-9001.05 | 2350 | 20.8 | 0.93i | -100 | J |
| k-1.1-9022.04 | 2500 | 10.2 | - | - | - |
| +1.1-3021.00 | 2600 | 3.5 | 2.0.3 | C. 380 | D |
| +1.1-9022.05 | 2000 | 3.7 | 2.032 | 0.3 cc | $D$ |
| \$-1.i-goci. 11 | zaćle | 0.50 | :3.08 | - 3 V20 | - |



| Station | Ground Range | Maximum Overpressure | Arrival <br> Time | Positive <br> Duration | Type Waveform |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | feet | psi | sec | sec |  |
| Shot Orens |  |  |  |  |  |
| 9-1.1-9021.014 | 500 | 225.0 | - | 0.255 | $\boldsymbol{\lambda}$ |
| 9-1.1-9021.01B | 600 | 115.- | - | - | - |
| 9-1.1-9021.01B | 600 | 102.5 | - | 0.148 | $A$ |
| 9-1.1-9021.01C | 800 | 62.0 | - | 0.228 | A |
| 9-1.1-9021.02 | 1000 | 37.0 | - | - | - |
| 9-1.1-9021.02 | 1000 | 35.5 | - | 0.282 | A |
| 9-1.1-9022.01 | 1150 | 27.0 | - | 0.354 | $A$ |
| 9-1.1-9021.03A | 1420 | 13.3 | - | 0.270 | B |
| 9-1.1-9022.02 | 1511 | 11.0 | - | 0.441 | B |
| 9-1.1-9022.02A | 1575 | 21.6 | - | 0.423 | B |
| 9-1.1-9022.02B | 1650 | 8.6 | - | 0.460 | B |
| 9-1.1-9021.04 | 1700 | 8.7 | - | 0.456 | B |
| 9-1.1-9022.03 | 1800 | 9.1 | - | 0.450 | B |
| 9-1.1-9022.03 | 1800 | 9.4 | - | - | - |
| 9-1.1-9021.08 | 3000 | 5.9 | - | - $\overline{8} 10$ | - |
| 9-i.1-9021.11 | 6600 | 2.3 | - | 0.840 | D |
| Shot Shasta |  |  |  |  |  |
| c-1.1-9020.01 | 1000 | 55.0 | - | - 707 | - |
| 2-1.1-9020.02 | 1500 | 25.8 | 0.231 | 0.407 | D |
| 2-1.1-9020.03 | 2000 | 14.0 | 0.578 | 0.448 | D |
| 2-1.1-9020.04 | 2500 | 10.0 | - | 0.559 | D |
| 2-1.1-9020.05 | 3000 | 8.0 | 1.165 | 0.622 | D |
| 2-1.1-9020.06 | 18,000 | 1.2 | 13.51 | 1.218 | D |
| Shot Gallleo |  |  |  |  |  |
| 1-1.1-9001.00 | 600 | 160.0 | - | - | - |
| 1-1.1-9001.01 | 1050 | 36.5 | - | 0.255 | $\boldsymbol{A}$ |
| 1-1.1-9002.01 | 1200 | 25.0 | - | 0.315 | A |
| 1-1.1-9001.02 | 2400 | 23.7 | - | 0.324 | $A$ |
| 1-1.1-9002.02 | $=700$ | 16.5 | - | 0.402 | $A$ |
| 1-1.1-9001.03 | 1950 | 12.2 | - | 0.436 | B |
| 1-1.1-9002.03 | 2100 | 10.5 | - | 0.463 | B |
| 1-1.1-9001.04 | 2200 | 10.7 | - | 0.476 | B |
| 1-1.1-9001.05 | 2350 | 9.8 | - | 0.510 | D |
| 1-1.1-9002.04 | 2500 | 10.0 | - | 0.475 | D |
| 1-1.1-9001.06 | 2600 | 9.2 | - | 0.503 | D |
| 1-1.1-9002.05 | 2765 | 8.3 | - | 0.581 | D |
| 1-1.1-9002.06 | 4700 | 4.2 | - | 0.745 | D |
| Shot Charleston |  |  |  |  |  |
| 9-1.9-9020.01 | 1100 | 17.5 | - | 0.410 | D |
| 9-1.9-9020.01 | 1100 | 17.2 | - | 0.332 | D |
| 9-1.9-9020.02 | 1300 | - | - | - | - |
| 9-1.9-9020.03 | 1500 | 14.5 | - | 0.433 | D |
| 9-1.9-9020.03 | 1500 | 15.0 | - | $0 \cdot 5$ | - |
| 9-1.9-9020.04 | 2300 | 12.2 | - | 0.520 | D |
| 9-1.9-9020.04 | 2300 | 13.0 | - | 0.490 | D |
| 9-1.1-9021.08 | 3000 | 8.5 | - | - | - |
| 9-1.1-9021.08 | 3000 | 8.4 | - | 0.577 | D |
| 9-1.1-9022.08 | 5500 | 3.8 | - | - | - |
| 9-1.1-9022.08 | 5500 | 3.7 | - | 0.737 | D |
| 9-1.1-9022.11 | 8000 | 2.1 | - | 1.060 | D |
| 9-1.1-9022.12 | 16,000 | 1.10 | - | - | - |
| 9-1.1-9022.13 | 21,120 | - | - | - | - |
| 9-1.1-9022.14 | 25,873 | 0.95 | - | - | - |
| 9-1.1-9022.14 | 25,873 | 0.85 | - | 1.250 | - |
| 9-1.1-9022.16 | 51,216 | 0.22 | - | - | - |
| 9-1.1-9022.17 | 73,000 | 0.07 | - | 1.580 | - |
| Shot Morgan |  |  |  |  |  |
| 9-1.1-9021.01 | 600 | 95.0 | - | - | A |
| 9-1.1-9021.02 | 1000 | 33.0 | - | 0.281 | $A$ |
| 9-1.1-9021.04 | 1250 | 20.0 | - | 0.397 | A |
| 9-1.1-9021.07 | 2200 | 7.7 | - | 0.461 | C |
| 9-1.1-9021.08 | 3000 | 5.6 | - | 0.574 | D |
| 9-1.1-9021.10 | 5000 | 2.8 | - | 0.684 | D |
| 9-1.1-9021.11 | 6600 | 2.1 | - | 0.829 | D |
| 9-1.1-9022.11 | 8000 | 1.45 | - | 0.844 | D |
| 9-1.1-9021.12 | 15,000 | 1.10 | - | - | - |
| 9-1.1-9022.13 | 21,120 | 0.90 | - | - | D |
| 9-1.1-9022.14 | 25,872 | 0.70 | - | 0.987 | D |
| 9-1.1-9022.15 | 25,000 | 0.32 | - | 1.004 | D |

TABLE 2.7 q-GACE RESULTS, 3-FCOT LEVEL, MAXIMUM VAWES

| Station | Grewn Range | rotal Pruesure | 8tatic Overperesure | Ereague Differenct $\left(P_{p}-P_{s}\right)$ | $\begin{aligned} & \text { Dyomide } \\ & \text { Presime } \\ & \text { q. } \end{aligned}$ | $\begin{gathered} \text { Mach } \\ \text { nubuer } \\ \text { u/a } \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | reet | pei | pel | pil | Pir |  |
| Shot Frentisn |  |  |  |  |  |  |
| 3-1.1-9003.01 | 800 | 10.6 | 8.1 | 2.5 | - | - |
| 3-1.1-9003.02 | 900 | 8.8 | 7.4 | 1.4 | - | - |
| 3-1.1-9003.03 | 1200 | 6.0 | 5.4 | 0.6 | - | - |
| 3-1.1-9003.04 | 1500 | 4.0 | 3.6 | 0.4 | - | - |
| Shot Wileon |  |  |  |  |  |  |
| 9-1.1-9021.01 | 600 | - | - | - | - | - |
| 9-1.1-9021.02 | 1000 | 205.0 | 35.5 | 196.0 | 120.0 | 2.5 |
| 9-1.1-9021.023 | 1000 | - | - | - | 4,0 | , |
| 9-1.1-9021.03 | 1250 | 84.0 | 24.5 | 69.0 | 44.0 | 1.6 |
| 9-1.1-9021.04 | 1700 | 27.5 | 9.6 | 18.4 | 14.5 | 1.0 |
| 9-1.2-9021.0411 | 1700 | - | - | - | - | - |
| 9-1.1-9021.05 | 1950 | 23.0 | 9.6 | 13.4 | 11.0 | 0.83 |
| 9-1.1-9021.07 | 2200 | 12.3 | 9.4 | 3.2 | 3.0 | 0.44 |
| Shot Hood |  |  |  |  |  |  |
| 9-1.1-3021.04 | 1700 | 129.0 | 48.0 | 120.0 | 72.0 | 2.0 |
| 9-1.1-9021.07 | 2200 | -0 | $\stackrel{-}{8}$ | - | - | - 0 |
| 9-1.1-9021.08 | 3000 | 20.0 | 8.3 | 12.0 | 11.0 | 0.81 |
| 9-1.1-9021.09 | 4000 | 12.3 | 6.3 | 6.2 | 6.0 | 0.66 |
| 9-1.1-9021.10 | 5000 | 8.3 | 6.2 | 2.1 | 1.9 | 0.40 |
| 9-1.1-9021.11 | 6000 | 5.4 | 4.8 | 0.7 | 0.6 | 0.22 |
| Sbot Orums |  |  |  |  |  |  |
| 9-1.1-9021.01 | 600 | 620.0 | 105.0 | 590.0 | 340.0 | 3.5 |
| 9-1.1-9021.02 | 1000 | 144.0 | 35.0 | 126.0 | 73.0 | 1.6 |
| 9-1.1-9022.02 | 1500 | 39.0 | 11.0 | 34.5 | 24.0 | 1.4 |
| 9-1.1-9021.04 | 1700 | 25.9 | 8.8 | 9.1 | 7.4 | 0.87 |
| Shot Eepler |  |  |  |  |  |  |
| 4-1.1-9022.01 | 1200 | - | - | - | - | - |
| 4-1.1-9022.02 | 1700 | $\stackrel{\square}{6}$ | - | - | - | - |
| 4-1.1-9022.03 | 2100 | 16.8 | 13.7 | 4.7 | 4.3 | 0.50 |
| 4-1.1-9022.04 | 2500 | 13.5 | 21.0 | 2.6 | 2.3 | 0.38 |
| 4-1.1-9082.05 | 2800 | 10.0 | 8.7 | 1.5 | 1.3 | 0.30 |
| Shot Gallleo |  |  |  |  |  |  |
| 9002-01 | 1200 | 55.0 | 24.5 | う".0 | 24.0 | 0.98 |
| 9002.02 | 1700 | 57.0 | 26.0 | 20.5 | 16.0 | 0.90 |
| 9002.03 | 2100 | 22.0 | 20.3 | 14.1 | 11.3 | 0.86 |
| 9002.05 | 2765 | 10.2 | 8.3 | 2.1 | 2.1 | 0.36 |




Figure 2.3 Wave shape versus ground distance for various yields.





$10^{3}$
Ground Range, Faet
Figure 2.10 Fatio of precursor maximum pressure to main shock peak pressure $P_{m}$ versus ground distance, scaled to 1 kt .
$\omega_{d} /{ }^{I_{d}} 0110 \mathrm{~d}$

эasw * $\omega_{\perp \nabla}$ 2w! 1

Figure 2.8 Time between arrival of
versus ground distance, scaled to 1 kt.
oəsw ${ }^{1} \downarrow \nabla$ əw! 1


Figure 2.11 Ratio of pressure $P_{p_{2}}$ to $P_{p_{1}}$
versus ground distance, scaled to 1 kt .



Figure 2.13 Predicted wave shapes, Shot Priscilla.

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Figure 2.29 Overpressure-time -orles, Shut Hood, at
distances of $2,400,3,000,4,590,0,000$, and 5,500 feet.

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Figure 2.32 Overpressure-time histories, Shot Kepler, at distances of $1,400,1,700,1,950,2,100$, and 2,200 feet.

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Figure 2.36 Overpressure-time histories, Shot Shasta, at distances of $1,500,2,000,2,500,3,000$, and 18,000 feet.

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Figure 2.43 Orerpressure-time histories, Shot Morgan, at distances of $6,600,8,000,25,872$, and 35,000 feet.








Figure 2.46 Pressure and Mach number versus time, at 1,700 feet, Shot Wilson.


Figure 2.47 Pressure and Mach number versus time, at 1,950 feet, Shot Wilson.







Time, msec

Figure 2.48 Pressure and Mach number versus time, at 2,200 feet, Shot Wilson.


Figure 2.49 Pressure and Mach number versus time, at 3,000 feet, Shot Hood.



Figure 2.50 Pressure and Mach number versus time, at 4,000 feet, Shot Hood.






Figure 2.51 Pressure and Mach number versus time, at 5,000 feet, Shot Hood.






Figure 2.52 Pressure and Mach number versus time, at 6,500 feet, Shot Hood.






Figure 2.53 Pressure and Mach number versus time, at 2,100 feet, Shot Kepler.






Figure 2.54 Pressure and Mach number versus time,

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Figure 2.55 Pressure and Mach number versus time, at 2,800 feet, Shot Kepler.








Figure 2.56 Pressure and Mach number versus time, at 1,000 feet, Shot Owens.





Figure 2.61 Maximum overpressure versus ground range, Shot $W^{i}$ ison.


Figure 2.62 Maxim'ım dynar:ic pressure versus ground range, Shot Wilsun.


Figure 2.63 Maximum overpressure versus ground range, Shot Hood.


Figure 2.64 Maximum dynamic pressure versus ground range, Shot Hood.


Figure 2.65 Maximum overpressure versus ground range, Shot Owens.

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Figure 2.66 Maximum dynamic pressure versus ground range, Shct Owens.
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Figure 2.67 Maximum overpressure versus ground range, Shot Kepler.

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Figure 2.68 Maximum dynamic pressure versus ground range, Shot Kepler.


Figure 2.69 Maximum overpressure versus ground range, Shot Shasta.



Figure 2.71 Maximum dynamic pressure versus ground range, Shot Galileo.



Figure 2.72 Maximum overpressure versus ground range, Shot Charleston.


Figure 2.73 Maximum overpressure versus ground range, Shot Morgan.



Figure 2.74 Arrival time and positive duration versus ground range, Shot Franklin.


Figure 2．75 Arrival time and positive duration versus ground range，Shot Wilson．


Figure 2.76 Arrival time and positive duration versus ground range, Shot Hood.


Figure 2.77 Positive duration versus ground range, Shot Owens.


Figure 2.78 Arrival time and positive duration versus ground range, Shot Kepler.


Figure 2.79 Arrival time and positive duration versus ground range, Shot Shasta.


Figure 2.80 Positive duration versus ground range, Shot Galileo.


Figure 2.81 Positive duration versus ground range, Shot Charleston.


Figure 2.82 Positive duration versus ground range, Shot Morgan.


Figure 2.83 Layout of shield configuration, Shot Kepler.



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Figure 2.86 Maximum avorpressure-distance versus ground range, scaled to 1 kt , Shot Fanklin.



Figure 2.87 Maximum overpressure-distance versus ground range, scaled to 1 kt , Shot Wilson.


Figure 2.89 Maximum overpressure-distance versus ground range, scaled to 1 kt , Sbot Owens.

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Figure 2.92 Maximum overpressure-distance versus
ground range, scaled to 1 kt , Shot Galileo.




Figure 2.94 Maximum overpressure-distance versus ground range, scaled to 1 kt , Shot Morgan.

Figure 2.95 Height-of-burst curves versus ground range for isopressure contours.

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Figure 2.96 Comparison of overpressure, Shots Wilson and Kepler.




Figure 2．97 Comparison of dynamic pressure versus ground range，Shots Wilson and Kepler．

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Figure 2.98 Comparison of arrival time versus ground range, Shots Wilson and Kepler.


Figure 2.99 Comparison of positive duration versus ground range, Shots Wilson and Kepler.



Figure 2.100 Comparison of overpressure from Shots Kepler, Shasta, and Galileo with the good surface height-of-burst curve of TM 23-200.

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Figure 2.101 Wave types from Plumbbob versus ground distance and yield.


Figure 2.102 Comparison of predicted and measured waveforms, Shot priscilla.


## Chapter 3

## BLAST PRESSURE, SHOT JOHN

Measurements of pressure waves at and near ground level resulting from nuclear bursts at high altitude can be used to determine proper scaling methods for overpressure, positive and negative phase durations, and time of arrival. Data from only one shot had been available for checking the procedures for high-altitude bursts. Additional data was required to substantiate the conclusions reached from these limited measurements. Shot John, burst at high altitude, provided some meag:rements.

### 3.1 OBJECTIVE

The objective of Project 1.1 was to obtain the overpressure-time history at various positions near and on the ground surface during Shot John. The data obtained was to corroborate further the scaling methods for free-air pressure and to determine the surface reflection of shock waves. The pressure measurements of interest were below the 1-psi region.

### 3.2 BaCKGROUND AND THEORY

From the studies conducted (Reference 18) of a nuclear weapon burst at high altitude, it was shown that the blast efficiency of the detonation will decrease as the altitude of burst increases. A check of this phenomenon would require a blast line located at the same elevation as the burst point in order that the characteristics of shock-wave propagation as a function of distance could be determined through a homogeneous atmosphere. A pressuredistance curve obtained by this means could be compared to a free-air curve obtained from a burst at sea-level conditions to give an indication of the blast efficiency of the device.

As pointed out in Reference 19, measurements of pressure waves at and near the ground surface from a high-altitude shot could be used to check on the scaling laws, provided that the yield, blast efficiency, and frec-air curve are known. The measurements taken during Operation Teapot (Reference 19) indicate that the modified Sachs scaling given in Equations 3.1, 3.2, and 3.3 is nore appropriate for scaling overpressures, whereas the ordinary Sachs scaling is more appropriate for scaling the time parameters of the pressure wave.

To scale overpressures by modified Sachs scaling (Reference 20) from ambient pressure conditions at which measurements were taken to sea-level conditions, the following apply:

$$
\begin{equation*}
P_{s o}=P_{s a}\left(\frac{P_{0}}{P_{a}}\right) \tag{3.1}
\end{equation*}
$$

and the distance scaled to 1 kt at sea level is:

$$
\begin{equation*}
R_{0}=R_{2}\left(\frac{P_{a}}{P_{0} W}\right)^{1 / 3} \tag{3.2}
\end{equation*}
$$

time duration of the blast wave is scaled in accordance with the following:

$$
\begin{equation*}
t_{0}=t_{2}\left(\frac{P_{a}}{P_{0} W}\right)^{1 / 3}\left(\frac{T_{a}}{T_{0}}\right)^{1 / 2} \tag{3.3}
\end{equation*}
$$

$$
\text { Where: } \begin{aligned}
\mathbf{P}_{\mathbf{S}} & =\text { overpressure } \\
\mathbf{P} & =\text { ambient pressure } \\
\mathbf{R} & =\text { slant distance } \\
\mathbf{W} & =\text { yield of weapon } \\
\mathbf{t} & =\text { time, duration of blast wave } \\
\mathbf{T} & =\text { absolute temperature }
\end{aligned}
$$

and the subscripts a and o refer to conditions at point of measurement and at sea level, respectively.

The ordinary Sachs scaling (Reference 21) for overpressure and time parameters of the shock wave take the same form, but $\mathbf{P}_{\mathbf{2}}$ and $\mathrm{T}_{\mathbf{a}}$ now are the ambient conditions at burst height.

### 3.3 PROCEDURE

The gages employed for measurement of the blast wave were self-recording very-lowpressure (VLP) gages (Reference 4). Basically, the VLP gage contains a diaphragm with a stylus attached to the center and a glass disk on which the motion of the diaphragm resulting from the pressure acting thereon is scribed by the stylus. The glass disk is on a turntable driven by a chronometrically governed motor. The initiation of the driving mechanism is accomplished at zero time (time of detonation) by photoelectric means.

Ten VLP gages were used to obtain the desired information. The gages were placed atop five poles 50 feet high and also at surface level. To check on the symmetry of the blast wave front, the gages were placed at ground zero and 10,000 feet from ground zero in four quadrants-north, south, east, and west (Figure 3.1).

### 3.4 DATA REDUCTION

The procedure for data reduction of the VLP records was similar to that used for $\mathbf{P}_{\mathbf{t}}{ }^{-}$ gage records. However, a correction was applied to the VLP records that showed pressure values higher than those reported previously-

Calibration of the VLP gage in the past yielded a curve that showed deflecticn versus applied pressure. The pressure to the gage was statically applied with enough time lapse between step increases for temperature equilibrium to be reached. The air inside the scaled chamber of the gage then obeys Boyles law, i.e., PV = coastant. For the case when the pressure applied is sudden, such as a shock wave, the gas inside the scaled chamber obeys the adiabatic compression law, $\mathrm{PV}^{\gamma}=$ constant. Furthermore, since the ratio of the volume of the scaled chamber to the area of the diaphragm is not large, the pressure inside the chamber was increased appreciably. These factors required that a new technique be devised for calibration of the VLP gage. A complete description of the calibration technique is given in Reference 22. The correction factor to be applied to the past data was determined to be $1 \times 10^{-6} \mathrm{psi} /$ mil deflection.

### 3.5 RESULTS

Of the 10 VLP gages used, good pressure-time records were obtained on 7, peak
pressures only were obtained on 2, and 1 failed to yield any information. The pressure measurements are given in Table 3.1. Traces of the pressure-time histories are given in Figures 3.2 through 3.5. The oscillations in the initial portion of each wave are due to lack of proper amount of damping. Dotted lines were drawn through these oscillations by averaging the top and bottom of each cycle. The maximum pressure was read from the smoothed curve. No attempt was made to determine the time of arrival or the positive duration of the blast wave.
3.6 DISCUSSION

The reflection factors derived from the pressure measurements are listed in Table 3.2. These were determined by calculating the ratio of pressure $\left(P_{4}-P_{3}\right)$ to $P_{2}$. The $P_{1}$ point is shown in Figure 3.6. The average reflection factor was 0.90 .

The reflected pressure $P_{R}$ measured from the surface gage was converted to incident pressure $P_{1}$ by use of the relation $P_{1}=P_{R} / 1.90$. The value 1.90 was based on the a ar age measured reflection factor of 0.90 .

Table 3.3 lists the scaling factors used for the incident pressures at 50 feet and the ground. The scaled values of pressure and slant range, using the modified Sachs scaling and standard Sachs scaling, are given in Table 3.4. These scaled values are plotted in Figure 3.7 along with the scaled free-air curve from TM 23-200 (Reference 7). From the compar' on of the measured scaled values with the free-air curve, the modified Sachs scaled values approach closer to the predicted frr . "curve than the standard Sachs scaled values. However, the modified Sachs sualed values themselves do not fit this curve too well.

The predicted free-air curve was constructed based on data derived in Shot HA of Operation Teapot (Reference 19). Curves for Shots HA and John are shown in Figure 3.8. For the Shot HA points, only the SC daid was used. In both methods of scaling the pressure, a considerable amount of spread in the data exists. The curves in Figure 3.8 do not indicate which scaling method is best suited for treatment of pressure data. To corrolate eitner of the scaling methods, additional information is needed. In particular, the information required is the partitioning of energy of detonations at various altitude-nr a free-air pressure versus distance curve at co-altitude.
3.7 CONCLUSIONS AND RECOMMENDATIONS

During Shot John, the instrumentation was such that the scaling methodis, Sachs or modified Sachs, were neither corroborated nor invalidated. For high-altitude bursts, pressure measurements should be made at co-altitude and lower altitudes for studying the partitioning of energy and propagation of shock waves through a nonhomogeneous atmosphere.

Unimanary
table 3.1 pressure measurements from vlp gages, shot john

| Ground Distance | Station | Position | $\begin{aligned} & \text { Altitude } \\ & \text { of } \\ & \text { Station } \end{aligned}$ | Slant <br> Range | $\begin{gathered} \text { Gage } \\ \text { No. } \end{gathered}$ | $\overline{P_{1}}$ Over- pressure | $\mathrm{P}_{2}$ Incident Over- pressure | $\mathrm{P}_{3}$ <br> Minimum <br> Pressure Before Reflection | $\mathrm{P}_{4}$ <br> Reflected Pressure |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| feet |  |  | feet | feet |  | psi | psi | psi | psi |
| 774 | Intended GZ | Ground 50-1t pole | 4,280 | 14,830 | $\begin{aligned} & \mathbf{1}^{*} \\ & 5^{*} \end{aligned}$ | $\begin{aligned} & 0.450 \\ & 0.470 \end{aligned}$ | $0.240$ | $0.120$ | 0.350 |
| 10.760 | North | Ground 50-ft pole. | 4,400 | 18,220 | $\begin{aligned} & 4^{*} \\ & 8^{*} \end{aligned}$ | $\begin{aligned} & 0.410 \\ & 0.280 \end{aligned}$ | $\overline{0.150}$ | $0.080$ | 0.210 |
| 9,910 | East | Grouni 50-ft pole | 4,520 | 17,640 | $\begin{aligned} & 6 \dagger \\ & 9^{*} \end{aligned}$ | $-$ | $\overline{0.230}$ | $-\overline{0.100}$ | 0.320 |
| 9,230 | South | Ground 50-ft pole | 4,240 | 17,550 | $\begin{aligned} & 3 \ddagger \\ & 2 \ddagger \end{aligned}$ | 0.280 | - | - | - |
| 10,450 | West | Ground 50-ft pole | 4,440 | 18,000 | $\begin{gathered} 10^{*} \\ 7^{*} \end{gathered}$ | $\begin{aligned} & 0.420 \\ & 0.430 \end{aligned}$ | $\overline{-}$ | 0.090 | 0.280 |


| Station | $\begin{aligned} & \text { Gage } \\ & \text { Number } \end{aligned}$ | Incident Pressure | $\mathrm{P}_{4}-\mathrm{P}_{3}$ | $\begin{aligned} & \text { Reflection } \\ & \text { Factor } \\ & \hline \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: |
| Intended GZ | 5 | 0.240 | 0.230 | 0.96 |
| North | 8 | 0.150 | 0.130 | 0.87 |
| East | 9 | 0.230 | 0.220 | 0.97 |
| West | 7 | 0.240 | 0.190 | 0.79 |
| Average |  |  |  | 0.90 |

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TABLE 3.3 SCALING FACTORS

| Station | Amblent Pressure， mb | Modified |  | $\begin{aligned} & \text { Sachs } \\ & S_{p} \end{aligned}$ | Scaling $S_{d}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $\begin{aligned} & \text { Sachs } \\ & s_{p} \end{aligned}$ | $\begin{gathered} \text { Scaling } \\ s_{d} \end{gathered}$ |  |  |
| Height of Burst | 502 | － | － | 2.018 | 0.6631 |
| Intended GZ | 870 | 1.167 | 0.7961 | － | － |
| North | 938 | 1.172 | 0.7944 | － | － |
| East | 860 | 1.178 | 0.7934 | － | － |
| South | 809 | 1.166 | 0.7959 | － | － |
| West | 863 | 1.174 | 0.7942 | － | － |

TABLE 3.4 SCALED PRESSURE－DISTANCE DATA，SHOT JOHN

| Station | Gage <br> No． | Slant <br> Range | Incident <br> Pressure | Modified <br> Slant <br> Range | Incident <br> Pressure |  | Sachs Scaling  <br> Slant  <br> Range  | Incident <br> Pressure |
| :--- | ---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | feet | psi | feet | psi | feet | psi |  |
| Intended |  |  |  |  |  |  |  |  |
| Ground | 1 | 14,830 | 0.237 | 11,810 | 0.277 | 9,830 | 0.478 |  |
| Zero | 5 | 14,790 | 0.240 | 11,770 | 0.280 | 9,810 | 0.484 |  |
| North | 4 | 18,220 | 0.216 | 14,470 | 0.253 | 12,080 | 0.436 |  |
|  | 8 | 18,150 | 0.150 | 14,420 | 0.176 | 12,040 | 0.303 |  |
| East | 9 | 17,600 | 0.230 | 13,960 | 0.271 | 11,670 | 0.464 |  |
| South | 3 | 17,550 | 0.147 | 13,970 | 0.171 | 11,640 | 0.297 |  |
| West | 10 | 18,000 | 0.221 | 14,300 | 0.259 | 11,940 | 0.446 |  |
|  | 7 | 17,560 | 0.240 | 13,950 | 0.282 | 11,640 | 0.484 |  |



Figure 3．：Field layout，Shot John．


Figure 3.4 Pressure-time history, VLP gage, east station, Shot John.


Figure 3.5 Pressure-time histories, VLP gages, west station, Shot John.


$$
\begin{aligned}
& \mathbf{P}_{1}=\text { reflected ground pressure } \\
& \mathbf{P}_{2}=\text { incident pressure at gage } \\
& \mathbf{P}_{3}=\text { pressure at gage before reflected pressure } \\
& \mathbf{P}_{4}=\text { reflected pressure at gage }
\end{aligned}
$$

Figure 3.6 Pressure-time trace, 50 feet above surface.



Figure 3.7 Comparison of maximum overpressure "ersus slant range, Shot John, with predicted curve, standard and modified Sachs scaling.


Figure 3.8 Maximum overpressure versus slant range, Shots John and HA, standard and modified Sachs scaling.

Chapter 4


## VLP-GAGE MEASUREMENTS


#### Abstract

At the beginning of Operation Plumbbob it became apparent, from the number of queries made concerning the predictions of pressure at large distances from the point of burst, that there was a paucity of data for which predictions of pressures could be made with a large degree of reliabilit;: Several agencies were concerned about the effects that the pressures at large distances from point of burst may have on equipment and structures. These pressures were measured to provide the desired information to interested agencies.


### 4.1 OBJECTIVE

The objective of this portion of Project 1.1 was to obtain the pressure-time history of the blast at large distances from several nuclear bursts.

### 4.2 PROCEDURE

VLP gages were used for measurement of pressure at large distances from point of burst.
The gages and the shots are given in Table 4.1. Some of the gages were used as backup to the instrumentation provided by Project 5.2.

These gages could be initiated either manually or by means of a photoelectric system. The standard method for initiation was manual and consisted of throwing a switch to complei the motor drive circuit and start the recording system. The photoinitiation system was usea at stations when it was impossible for personnel to man the gage.

### 4.3 RESULTS

The shots and the pressure measurements obtained are given in Table 4.1. The distances given in the table are approximate. These were obtained by taking measurements from a map of the area prepared by the Corps of Engineers.

The measured values for all shots were scaled to 1 kt , sea level, and plotted in Figure 4.1. A large spread of the data exists, and the dotted curves indicate the variation in pressure that can be expected. The lack of instrumentation between burst point and location of gage, for determining atmospheric conditions, did not warrant further analysis of the terrain effects on the propagation of the shock waves.

The pressure-time histories of the records are given in Figures 4.2 through 4.8.

### 4.4 CONCLUSION

A large variation in maximum overpressures at large distances can be expected from nuclear detonations.




Figure 4.1 Maximum overpressure versus slant range, VLP gages, scaled to 1 kt , modified Sachs.



Figure 4.4 Overpressure-time histories, VLP Gages 4 and 5, Shot Owens.



Figure 4.6 Overpressure-time histories, VLP Gages 6 and 9, Shot Stokes; VLP Gage 5, Shot Whitney.





Figure 4.7 Overpressure-time histories, VLP Gages 2, 8, and 10, Shot Charleston.


Figure 4.8 Overpressure-time histories, VLP Gages 2 and 10, Shot Morgan.

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[^0]:    - See foximote. Tahle 2.1

